

Lecture 35. Spectral Filters for LIDAR Applications

- Spatial filters
- -- Considerations on beam divergence and FOV
- Interference filters
- Daytime filters: Fabry-Perot interferometers
- Daytime filters: Faraday filters
- □ Atomic and molecular absorption filters

Spatial Filters Considerations on Beam Divergence & FOV

Small laser beam divergence → small receiver field of view (FOV)

-- A very effective way to reduce daytime background

Whether a small beam divergence thus small FOV can be used is mainly determined by two factors –

- 1) Will the atmosphere saturation be resulted? If yes, then we cannot use small divergence and FOV.
- 2) Is the laser spatial mode good enough (TEMOO mode) to achieve small beam divergence, enabling small FOV?

How much FOV is suitable? FOV is usually larger than the claimed beam divergence to tolerate the fluctuations in the beam pointing.





IF and FPI are based on multiple beam interference.

LIDAR REMOTE SENSING



Multiple Beam Interference



Phase difference between two adjacent beams is given by

$$\delta = \frac{2\pi}{\lambda} \Delta L = \frac{2\pi}{\lambda} 2nh\cos i$$

□ In the figure, t and r are the amplitude-transmission and reflection coefficients. Intensity transmission and reflectivity are the square of t and r, i.e., $T = t^2$ and $R = r^2$.



Multiple Beam Interference

Total amplitude is the sum of multiple beam amplitudes

$$\tilde{U}_T = Att'(1+r^2e^{i\delta}+r^4e^{2i\delta}+\ldots) = \frac{Att'}{1-r^2e^{i\delta}}$$

Thus, the transmission intensity is

$$I_T = \tilde{U}_T \tilde{U}_T^* = \frac{A^2 (tt')^2}{\left(1 - r^2 e^{i\delta}\right) \left(1 - r^2 e^{-i\delta}\right)} = \frac{I_0 (1 - r^2)^2}{1 - 2r^2 \cos \delta + r^4}$$

Recall $R = r^2$, therefore, we have

$$I_T = \frac{I_0}{1 + \frac{4R\sin^2(\delta/2)}{(1-R)^2}}$$

Multiple Beam Interference

□ When $\delta = 2k\pi$, the transmission light reaches maximum, which determines the transmission wavelength or frequency.

□ If incident angle i = 0, transmission wavelengths and frequencies are determined by $C = \frac{1}{C} \frac{1}{k}$

$$2nh = k\lambda_k$$

$$v_k = \frac{c}{\lambda_k} = \frac{kc}{2nh}$$

Thus, the frequency spacing or Free-Spectral-Range is



 $\Delta v_k =$

Full-Width-at-Half-Maximum for each transmission line is

Finesse is defined as

$$F = \frac{FSR}{\Delta v_k} = \frac{\pi \sqrt{R}}{1 - R}$$



Interference Fringes



Periodic transmission lines



Fabry-Perot Etalon





Typical Transmittance of Fabry-Perot Etalon





Interference Filter





Spacer at half wavelength for the desired wavelength, or a multiple of that. □ Interference filters are multilayer thin-film devices, based on Fabry-Perot interferometer.

Constructive interference at the desired wavelength (spacer d = $\lambda_0/2$)

- round trip 2d = λ_0

Destructive interference at other wavelengths to block them.



Multi-Cavity Interference Filter

Multi-cavity interference filter gives better rejection to wavelengths other than desired wavelength



Note: The actual FWHM will be different in each case.

Effect of number of cavities on passband shape (normalized transmittance for ZnS/Na₃AIF₆ passband interference filters with 10-nm FWHM)



Interference Filter



Colored glass etc is optional to further block wavelengths far away from desired wavelength, especially at shorter end.



Combination of Interference Filter and Etalon in Lidar Receiver



□ Typical bandwidth of F-P etalon in lidar receiver is about 10–30 GHz.



Spectral Filters: Faraday Filters



FADOFs (Faraday Anomalous Dispersion Optical Filters) are based on the magnetic-field induced Faraday rotation effect of linear polarization of light.



Faraday Effect

□ Faraday effect is the rotation of light polarization by some media under magnetic field. β_{N}





Refraction index n of dilute Na vapor

$$n = \sqrt{1+\chi} \cong 1 + \frac{1}{2}\chi = 1 + \frac{1}{2}\chi' - i\frac{1}{2}\chi''$$
 (5.74)

 $\boldsymbol{\chi}$ is the electric susceptibility of Na vapor



Optical Activity vs. Waveplate Retardation



Optical activity: linearly polarized light propagates along the optic axis with the linear polarization rotating along the way. -- "Circular polarization birefringence"

Figure 1.14 Rotation of the plane of polarization by an optically active medium.



Waveplate retardation: linearly polarized light propagates perpendicular to the optic axis and the polarization changes its state along the way.

The decomposition of linearly polarized light is based on the orthogonal eigenstates of a crystal. For example, optical activity has L and R circular polarizations as its eigen-vibrations, while wave plate has o and e linear polarization as its eigne-vibrations. PROF. XINZHAO CHU CU-BOULDER, SPRING 2016



Faraday Rotator: Magnetic-Field-Induced Optical Activity



图 6-79 磁致旋光

Magneto-Optic (MO) crystal: Field-induced optical activity or magnetically induced optical activity (polarization rotation) Light must propagate along the optic axis which is defined by the external magnetic field. PROF. XINZHAO CHU CU-BOULDER, SPRING 2016



Faraday Rotator: Magnetic-Field-Induced Optical Activity



Naturally occurring optical activity

Faraday (B-fieldinduced) optical activity

图 6-80 自然旋光与磁致旋光的比较



The light speed in vacuum c is given by $_{\mathcal{C}}$

- ε_0 dielectric constant in vacuum (the permittivity of a vacuum)
- $\mu_{\rm 0}$ permeability in vacuum

The phase velocity of EM wave in medium v is given by

- ε relative dielectric constant
- μ relative permeability

$$v = \frac{1}{\sqrt{\varepsilon_0 \varepsilon \mu_0 \mu}} = \frac{c}{n}$$

The refraction index of the medium is given by $n=\sqrt{arepsilon\mu}$

Except for ferromagnetic materials, $\mu \approx 1$ is true for most materials, so

$$n = \sqrt{\varepsilon \mu} \approx \sqrt{\varepsilon} = \sqrt{1 + \chi}$$

X – macroscopic electric susceptibility of the medium $\varepsilon_m = \varepsilon_0 \varepsilon$ – the dielectric constant (permittivity) of the medium



Faraday Effect

□ Faraday effect is the rotation of light polarization by some media under magnetic field. β_{N}





Refraction index n of dilute Na vapor



$$n = \sqrt{1+\chi} \cong 1 + \frac{1}{2}\chi = 1 + \frac{1}{2}\chi' - \mathrm{i}\frac{1}{2}\chi'' \tag{5.74}$$

 $\boldsymbol{\chi}$ is the electric susceptibility of Na vapor

Faraday Effect under Zeeman Splitting





Resonance Absorption vs. Dispersion







Explanation of Optical Activity



When light propagates along the optic axis, it can be decomposed to the superposition (synthesis) of a left and a right circularly polarized light.

If two circular polarizations experience different refraction indexes, extra phase shift will be introduced.

Consequently, two circular polarizations recombine to form a linear polarization with its orientation rotated.

$$\varphi_{\rm L} = \frac{2\pi}{\lambda} n_{\rm L} d,$$
$$\varphi_{\rm R} = \frac{2\pi}{\lambda} n_{\rm R} d,$$

$$\psi = \frac{1}{2}(\varphi_{\mathrm{R}} - \varphi_{\mathrm{L}}) = \frac{\pi}{\lambda}(n_{\mathrm{R}} - n_{\mathrm{L}})d.$$

图 6-72 旋光性的解释



 $\Delta \varphi = 2\pi \frac{l\Delta n}{\lambda}$

Faraday Filter

- Phase shift between two circular polarizations
- \Box When the phase shift is π , the polarization is rotated by $\pi/2$





Narrowband Daytime Filter: Faraday Filter @ Different B and T



[Kiefer et al., Nature Scientific Reports, 2014]



Narrowband Daytime Filter: Faraday Filter @ Different B and T





Narrowband Daytime Filter: Faraday Filter @ Different B and T





Fluorescence and Rayleigh Scattering vs. Doppler Shift





Atomic and Molecular Absorption Filter



Fig. 5.5. Calibration scan showing the transmission of the molecular (blue) and combined (green) channels as a function of frequency. The Doppler broadened molecular spectrum for 300 K is also shown (black). Line 1109 of the iodine absorption spectrum (central notch) rejects most of the aerosol scattering and the central portion of the molecular scattering while passing the wings of the molecular line. The spectral transmission of the combined channel is determined by the pre-filter etalon.