

Lecture 11. Temperature Lidar (1) Overview and Doppler Technique

- Overview of Temperature Measurement Techniques
- Doppler, Boltzmann, Integration and Rotational Raman
- Doppler Technique to Measure Temperature and Wind
- Doppler Shift and Broadening in Resonance Absorption
- > Doppler Shift and Broadening in Resonance Fluorescence
- Doppler Shift and Broadening in Rayleigh Scattering
- Resonance Fluorescence Doppler vs Rayleigh Doppler
- ☐ Na Doppler Lidar: Scanning vs. Ratio Techniques
- Summary



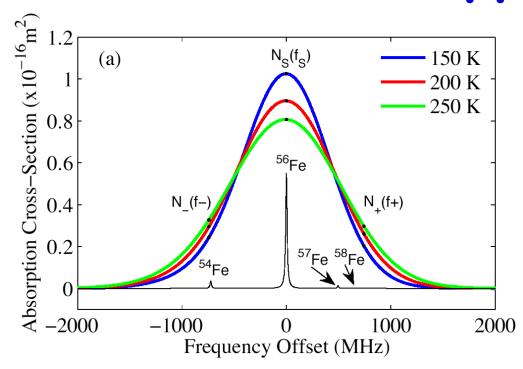
Temperature Measurement Techniques

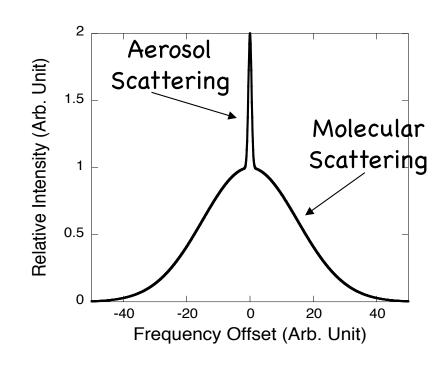
Use temperature-dependent effects or phenomena

- Doppler Technique Doppler broadening (not only for Na, K, and Fe, but also for Rayleigh scattering, as long as Doppler broadening dominate and can be detected)
- □ Boltzmann Technique Boltzmann distribution of atomic/molecular populations on different energy levels (not only for Fe, but also for molecular spectroscopy in optical remote sensing)
- ☐ Integration Technique (Rayleigh or Raman) integration lidar technique using ideal gas law and assuming hydrostatic equilibrium (not only for modern lidar, but also for cw searchlight and rocket falling sphere - some way to measure atmosphere number density)
- Rotational Raman Technique temperature dependence of population ratio, similar to Boltzmann technique



Overview: Doppler Technique





Atomic Absorption Line

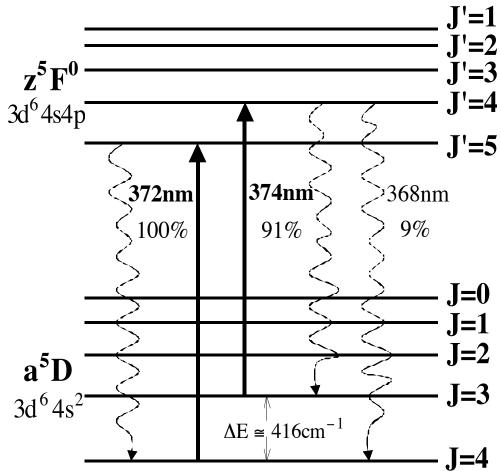
Rayleigh Scattering

$$\sigma_{rms} = \frac{v_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}} = \sqrt{\frac{k_B T}{M \lambda_0^2}}$$

$$\sigma_{rms} = 2v_0/c\sqrt{k_BT/M} = \frac{2}{\lambda_0}\sqrt{k_BT/M}$$



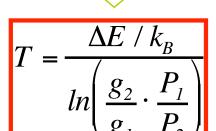
Overview: Boltzmann Technique



Atomic Fe Energy Level [Gelbwachs, 1994; Chu et al., 2002]

Maxwell-Boltzmann Distribution in Thermal-dynamic Equilibrium

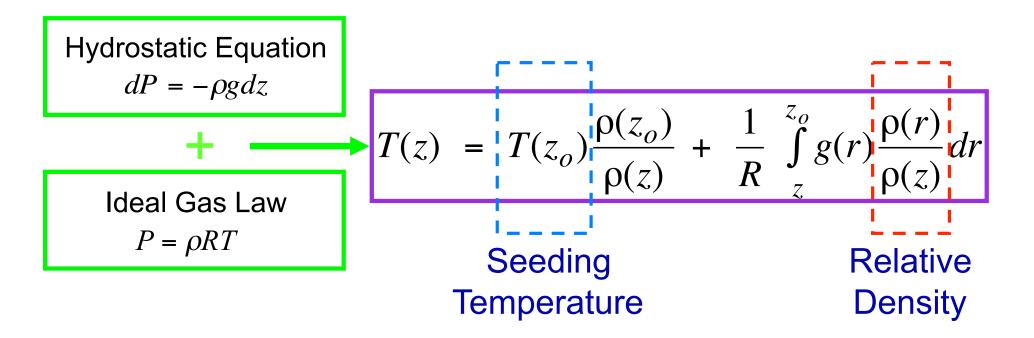
$$\frac{P_2(J=3)}{P_1(J=4)} = \frac{\rho_{Fe(374)}}{\rho_{Fe(372)}} = \frac{g_2}{g_1} \exp(-\Delta E/k_B T)$$



 P_1, P_2 -- Fe populations g_1, g_2 -- Degeneracy k_B -- Boltzmann constant T -- Temperature



Overview: Integration Technique



 $T(z_0)$ - Seeding Temperature;

R - gas constant for dry air;

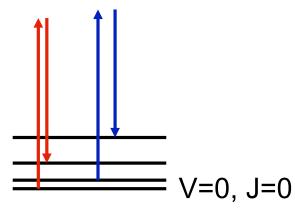
ρ - number density

g - gravitational acceleration

Number Density Ratio ⇒ **Temperature** (lidar backscatter ratio at different altitudes)

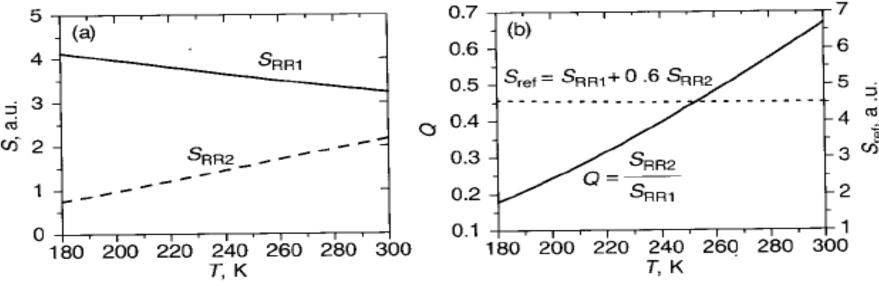


Overview: Rotation Raman Technique



$$Q(T) = \frac{\sum_{i=O_2,N_2} \sum_{J_i} \tau_{RR2}(J_i) \eta_i \left(\frac{d\sigma}{d\Omega}\right)_{\pi}^{RR,i} (J_i)}{\sum_{i=O_2,N_2} \sum_{J_i} \tau_{RR1}(J_i) \eta_i \left(\frac{d\sigma}{d\Omega}\right)_{\pi}^{RR,i} (J_i)}$$

V=0, J=0 where τ is the receiver transmission at RR line, η is the relative volume abundance of $N_2 \& O_2$.



Temperature can be derived from the ratio of two pure Rotational Raman line intensities. This is essentially the same principle as Boltzmann temperature technique!



Doppler Technique to Measure Temperature and Wind

- Doppler effect is commonly experienced by moving particles, such as atoms, molecules, and aerosols. It is the apparent frequency change of radiation that is perceived by the particles moving relative to the source of the radiation. This is called Doppler shift or Doppler frequency shift.
- □ Doppler frequency shift is proportional to the radial velocity along the line of sight (LOS) of the radiation –

$$\omega = \omega_0 - \vec{k} \cdot \vec{v}$$

$$\Delta \omega = \omega - \omega_0 = -\vec{k} \cdot \vec{v} = -\omega_0 (v/c) \cos \theta$$

$$\Delta v = -v_0 (v/c) \cos \theta = -(v/\lambda_0) \cos \theta$$

where ω_0 is the radiation frequency at rest, ω is the shifted frequency, k is the wave vector of the radiation (k=2 π / λ), and v is the particle velocity.



Doppler Technique

Due to particles' thermal motions in the atmosphere, the distribution of perceived frequencies for all particles mirrors their velocity distribution. According to the Maxwellian velocity distribution (Gaussian),

$$P(v_R \rightarrow v_R + dv_R) \propto \exp(-Mv_R^2/2k_BT)dv_R$$

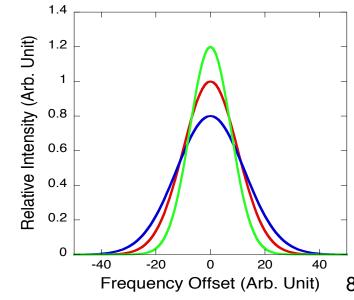
$$\omega = \omega_0 + \vec{k} \cdot \vec{v} = \omega_0 \left(1 + \frac{v_R}{c} \right) \longrightarrow v_R = \frac{\omega - \omega_0}{\omega_0 / c} = \frac{v - v_0}{v_0 / c}$$

Substituting v_R into the probability distribution, we obtain the power spectral density distribution (i.e., intensity versus the perceived frequency by moving particles) as a Gaussian lineshape,

$$I \propto \exp\left(-\frac{M(v-v_0)^2}{2k_BT(v_0/c)^2}\right)(c/v_0)dv$$

This is called Doppler broadening of a line. The peak is at $\omega = \omega_0$ and the rms width is

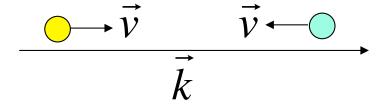
$$\sigma_{rms} = \frac{v_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}}$$





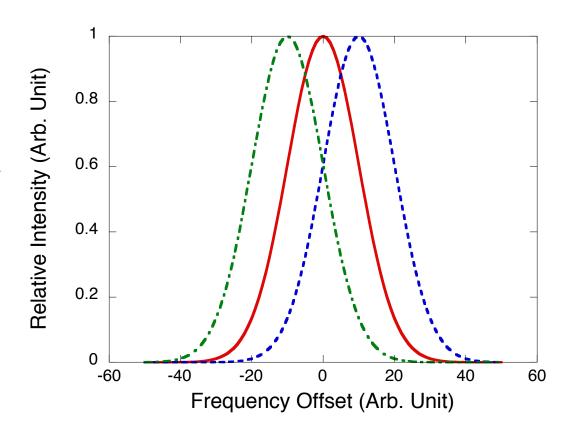
Doppler Shift in Resonance Absorption

$$\Delta\omega = \omega - \omega_0 = -\vec{k} \cdot \vec{v} = -\omega_0 \frac{v \cos \theta}{c}$$
 (12.13)



Emitter and receiver move towards each other:

- -Blue shift in perceived radiation frequency
- -Red shift in absorption peak frequency

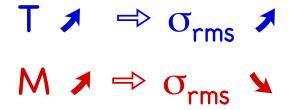


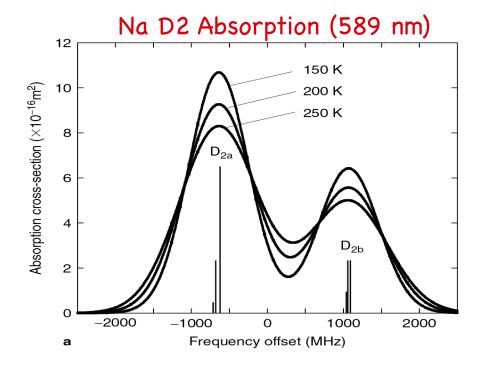
☐ The velocity measurements of lidar, radar, and sodar all base on the Doppler shift principle!

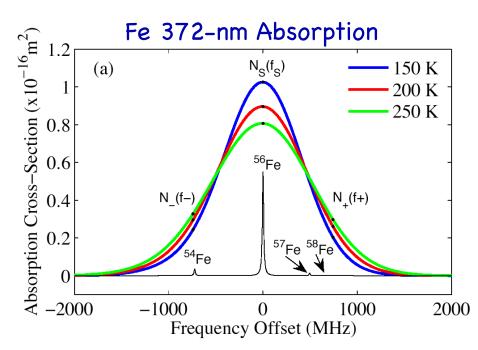


Doppler Broadening in Resonance Absorption Lines

$$\sigma_{rms} = \frac{v_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}}$$









Doppler Shift and Broadening in Resonance Fluorescence

☐ When an atom emits a resonance fluorescence photon, the photon has Doppler shift relative to the center freq. of the atomic absorption line as

$$\omega = \omega_0 + \vec{k} \cdot \vec{v} = \omega_0 \left(1 + \frac{v_R}{c} \right) \longrightarrow v_R = \frac{\omega - \omega_0}{\omega_0 / c} = \frac{v - v_0}{v_0 / c}$$

 \square According to the Maxwellian velocity distribution, the relative probability that an atom/molecule in a gas at temperature T has its velocity component along the line of sight between v_R and v_R+dv_R is

$$P(v_R \rightarrow v_R + dv_R) \propto \exp(-Mv_R^2/2k_BT)dv_R$$

 \square Substitute the v_R equation into the Maxwellian distribution,

$$I \propto \exp\left(-\frac{M(v-v_0)^2}{2k_BT(v_0/c)^2}\right)(c/v_0)dv$$

☐ Therefore, the rms width of the Doppler broadening is

$$\sigma_{rms} = v_0 / c \sqrt{k_B T / M} = \frac{1}{\lambda_0} \sqrt{k_B T / M}$$
 1 time



Doppler Shift in Rayleigh Scattering

Refer to textbook 5.2.2.4 Lidar wind vs radar wind measurements

Momentum Conservation

Energy Conservation

$$m\vec{v}_{1} + \hbar \vec{k}_{1} = m\vec{v}_{2} + \hbar \vec{k}_{2}$$

$$\frac{1}{2}mv_{1}^{2} + \hbar \omega_{1} = \frac{1}{2}mv_{2}^{2} + \hbar \omega_{2}$$

$$\omega_1 = \omega_2 + \vec{k}_1 \cdot \vec{v}_1 - \vec{k}_2 \cdot \vec{v}_2 + \frac{\hbar k_1^2}{2m} - \frac{\hbar k_2^2}{2m}$$

For Rayleigh or radar backscatter signals, we have

$$\vec{k}_2 \approx -\vec{k}_1$$
 $\vec{v}_2 \approx \vec{v}_1$

The frequency shift for Rayleigh or radar backscattering is

$$\Delta \omega_{Rayleigh,backscatter} = \omega_2 - \omega_1 = -2\vec{k}_1 \cdot \vec{v}_1$$



Doppler Broadening in Rayleigh Scatter

☐ To derive the Doppler broadening, let's write the Doppler shift as

$$\omega = \omega_0 \left(1 - \frac{2v_R}{c} \right) \longrightarrow v_R = \frac{\omega_0 - \omega}{2\omega_0 / c} = \frac{v_0 - v}{2v_0 / c}$$

 \square According to the Maxwellian velocity distribution, the relative probability that an atom/molecule in a gas at temperature T has its velocity component along the line of sight between v_R and v_R +d v_R is

$$P(v_R \rightarrow v_R + dv_R) \propto \exp(-Mv_R^2/2k_BT)dv_R$$

 $lue{}$ Substitute the v_R equation into the Maxwellian distribution,

$$I \propto \exp\left(-\frac{M(v_0 - v)^2}{2k_B T(2v_0/c)^2}\right) (c/2v_0) dv$$

☐ Therefore, the rms width of the Doppler broadening is

$$\sigma_{rms} = 2v_0/c\sqrt{k_BT/M} = \frac{2}{\lambda_0}\sqrt{k_BT/M}$$
 2 times !



Doppler Effect in Rayleigh Scattering

☐ In the atmosphere when aerosols present, the lidar returns contains a narrow spike near the laser frequency caused by aerosol scattering riding on a Doppler broadened molecular scattering profile.

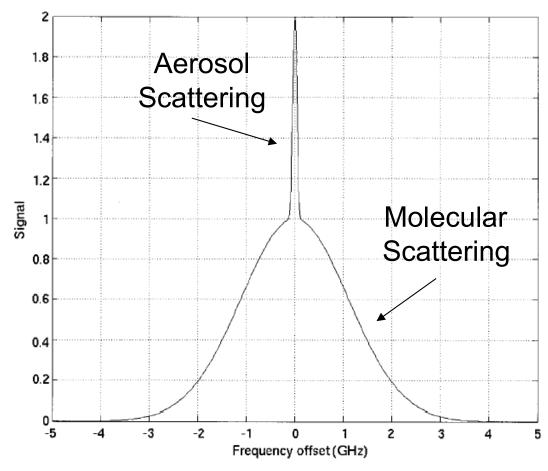


Fig. 5.1. Spectral profile of backscattering from a mixture of molecules and aerosols for a temperature of 300 K. The spectral width of the narrow aerosol return is normally determined by the line width of the transmitting laser.

At T = 300 K, the Doppler broadened FWHM for Rayleigh scattering is 2.58GHz, not 1.29GHz. Why?

Because Rayleigh backscatter signals have 2 times of Doppler shift!

Courtesy of Dr. Ed Eloranta
University of Wisconsin



Resonance Fluorescence Doppler versus Rayleigh Doppler

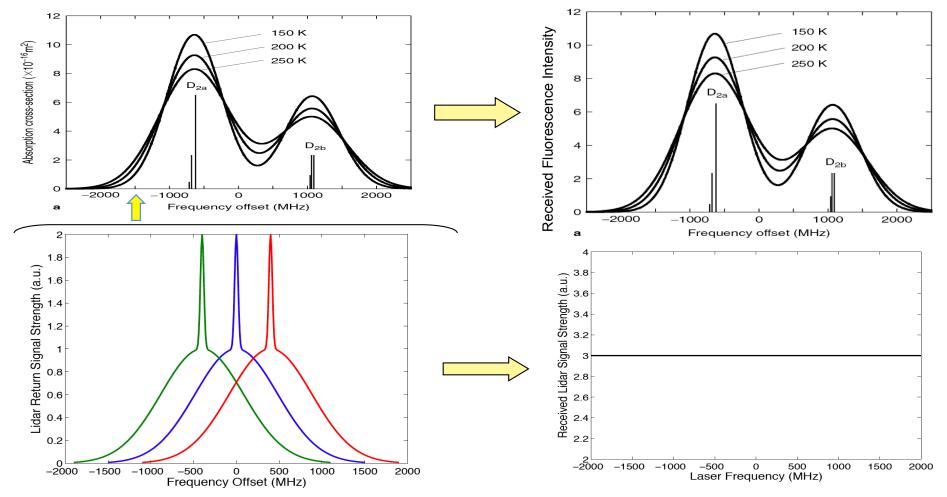
☐ Atomic absorption lines provide a natural frequency analyzer or frequency discrimination. This is because the absorption cross section undergoes Doppler shift and Doppler broadening. Thus, when a narrowband laser scans through the absorption lines, different absorption and fluorescence strength will be resulted at different laser frequencies. By using a broadband receiver to collect the returned resonance fluorescence, we can easily obtain the line shape of the absorption cross section so that we can infer wind and temperature. There is no need to measure the fluorescence spectrum. - Resonance fluorescence Doppler technique

☐ Rayleigh scattering also undergoes Doppler shift and broadening, however, it is not frequency discriminated. In other words, when scanning a laser frequency, the backscattered Rayleigh signal gives nearly the same Doppler broadened line width, independent of laser frequency. Thus, the atmosphere molecule scattering does not provide frequency discrimination. A frequency analyzer must be implemented into the lidar receiver to discriminate the return light frequency, i.e., analyze Rayleigh scattering spectrum to infer wind and temperature. - Rayleigh Doppler technique



Resonance Fluorescence Doppler versus Rayleigh Doppler

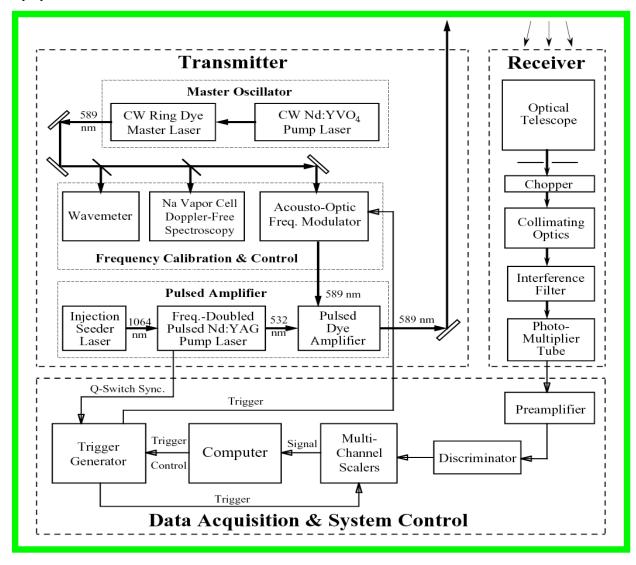
☐ How will the lidar return signal strength (vs. laser frequency) change when the lidar receiver is broadband and we scan the narrowband laser frequency - in resonance fluorescence case and in Rayleigh scattering case?





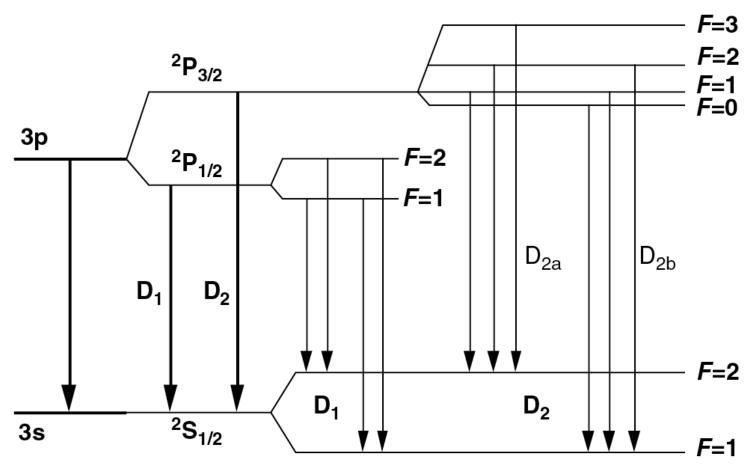
Na Doppler Wind and Temperature Lidar

☐ Na Doppler lidar is one of the most successful lidars.





Na Atomic Energy Levels

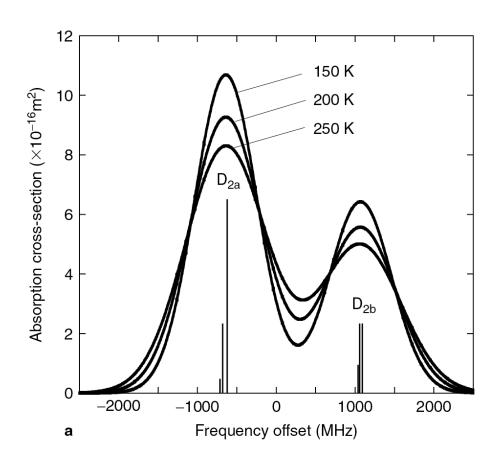


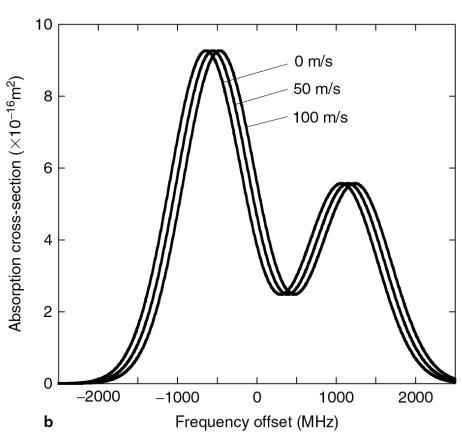
Na fine structure

Na hyperfine structure



Doppler Effect in Na D₂ Line Resonance Fluorescence





Na D₂ absorption linewidth is temperature dependent

Na D₂ absorption peak freq is wind dependent



Na Atomic Parameters

Table 5.1 Parameters of the Na D₁ and D₂ Transition Lines

Transition Line	Central Wavelength (nm)	$\begin{array}{c} \text{Transition} \\ \text{Probability} \\ (10^8\text{s}^{-1}) \end{array}$	Radiative Lifetime (nsec)	$egin{aligned} ext{Oscillator} \ ext{Strength} \ ext{$f_{ ext{ik}}$} \end{aligned}$
$\begin{array}{c} \hline D_1 \ (^2P_{1/2} {\to}^2S_{1/2}) \\ D_2 \ (^2P_{3/2} {\to}^2S_{1/2}) \end{array}$	589.7558 589.1583	0.614 0.616	16.29 16.23	0.320 0.641
Group	${}^{2}\mathrm{S}_{1/2}$	$^2\mathrm{P}_{3/2}$	Offset (GHz)	Relative Line Strength ^a
$\overline{\mathrm{D}_{2\mathrm{b}}}$	$F\!=\!1$	$F=2 \ F=1 \ F=0$	1.0911 1.0566 1.0408	5/32 5/32 2/32
$ m D_{2a}$	$F\!=\!2$	$egin{array}{c} F=0 \ F=3 \ F=2 \ F=1 \end{array}$	-0.6216 -0.6806 -0.7150	14/32 5/32 1/32
Doppler-Fre	ee Saturation–A	Absorption Fe	atures of the N	a D ₂ Line
, , .	•	f _b (MHz) 1067.8	$f_+(\mathrm{MHz}) = -21.4$	f_ (MHz) -1281.4

^aRelative line strengths are in the absence of a magnetic field or the spatial average. When Hanle effect is considered in the atmosphere, the relative line strengths will be modified depending on the geomagnetic field and the laser polarization.



Doppler-Limited Na Spectroscopy

 \blacksquare Doppler-broadened Na absorption cross-section is approximated as a Gaussian with rms width $\sigma_{\!\scriptscriptstyle D}$

$$\sigma_{abs}(\mathbf{v}) = \frac{1}{\sqrt{2\pi}\sigma_D} \frac{e^2 f}{4\epsilon_0 m_e c} \sum_{n=1}^{6} A_n \exp\left(-\frac{[\mathbf{v}_n - \mathbf{v}(1 - V_R/c)]^2}{2\sigma_D^2}\right)$$

- lacktriangle Assume the laser lineshape is a Gaussian with rms width $\sigma_{\!\scriptscriptstyle L}$
- ☐ The effective cross-section is the convolution of the atomic absorption cross-section and the laser lineshape

$$\sigma_{eff}(v) = \frac{1}{\sqrt{2\pi}\sigma_e} \frac{e^2 f}{4\epsilon_0 m_e c} \sum_{n=1}^{6} A_n \exp\left(-\frac{[v_n - v(1 - V_R/c)]^2}{2\sigma_e^2}\right)$$

where

$$\sigma_e = \sqrt{\sigma_D^2 + \sigma_L^2}$$
 and $\sigma_D = \sqrt{\frac{k_B T}{M \lambda_0^2}}$

The frequency discriminator/analyzer is in the atmosphere!



Doppler Scanning Technique

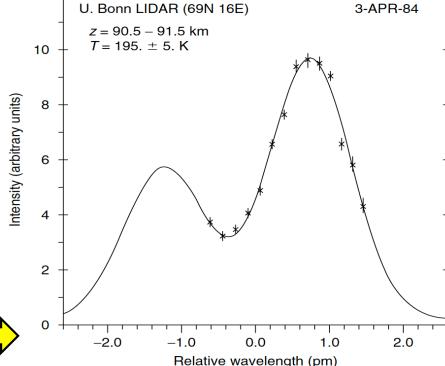
$$N_{Na}(\lambda,z) = \left(\frac{P_L(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{eff}(\lambda)n_{Na}(z)\Delta z\right) \left(\frac{A}{4\pi z^2}\right) \left(\eta(\lambda)T_a^2(\lambda)T_c^2(\lambda,z)G(z)\right)$$

$$N_{R}(\lambda, z_{R}) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{R}(\pi, \lambda)n_{R}(z_{R})\Delta z\right) \left(\frac{A}{z_{R}^{2}}\right) \left(\eta(\lambda)T_{a}^{2}(\lambda, z_{R})G(z_{R})\right)$$

$$\sigma_{eff}(\lambda, z) = \frac{C(z)}{T_c^2(\lambda, z)} \frac{N_{Na}(\lambda, z)}{N_R(\lambda, z_R)}$$

where
$$C(z) = \frac{\sigma_R(\pi, \lambda)n_R(z_R)}{n_{Na}(z)} \frac{4\pi z^2}{z_R^2}$$

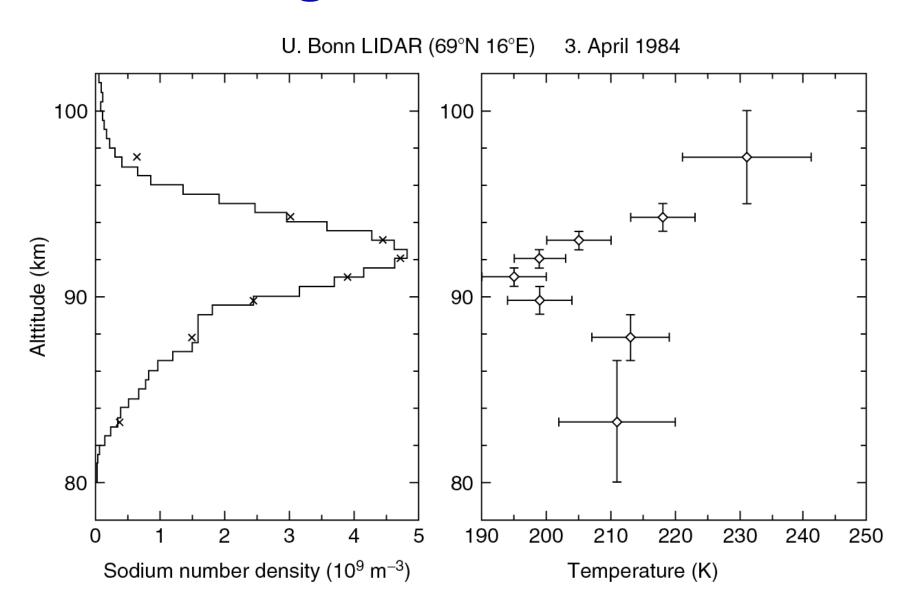
Least-square fitting gives temp [Fricke and von Zahn, JATP, 1985]





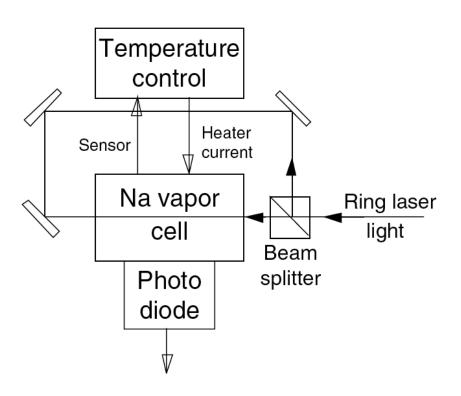


Scanning Na Lidar Results

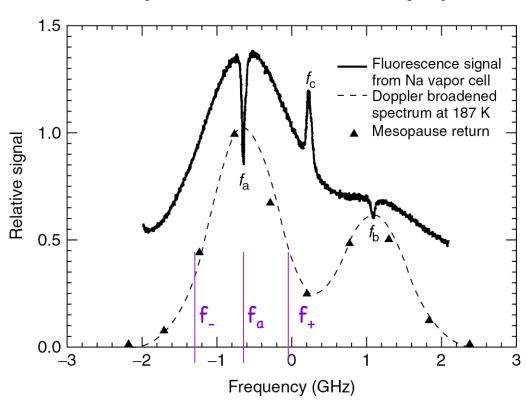


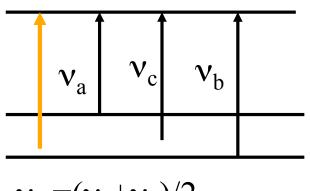


Doppler-Free Na Spectroscopy



See detailed explanation on Na Doppler-free saturation-fluorescence spectroscopy in Textbook Chapter 5.2.2.3.2



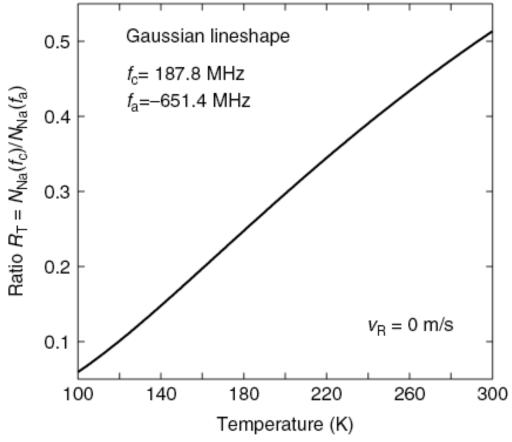


$$v_c = (v_a + v_b)/2$$



2-Frequency Doppler Ratio Technique

$$R_T(z) = \frac{N_{norm}(f_c, z, t_1)}{N_{norm}(f_a, z, t_2)} = \frac{\sigma_{eff}(f_c, z)n_{Na}(z, t_1)}{\sigma_{eff}(f_a, z)n_{Na}(z, t_2)} \approx \frac{\sigma_{eff}(f_c, z)}{\sigma_{eff}(f_a, z)}$$



$$N_{norm}(f,z,t) = \frac{N_{Na}(f,z,t)}{N_{R}(f,z,t)T_{c}^{2}(f,z)} \frac{z^{2}}{z_{R}^{2}}$$

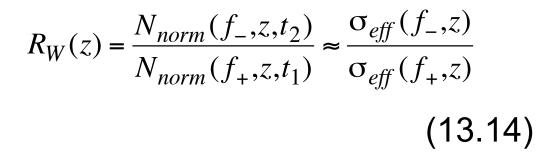
$$N_{norm}(f,z,t) = \frac{\sigma_{eff}(f)n_{Na}(z)}{\sigma_{R}(\pi,f)n_{R}(z_{R})} \frac{1}{4\pi}$$

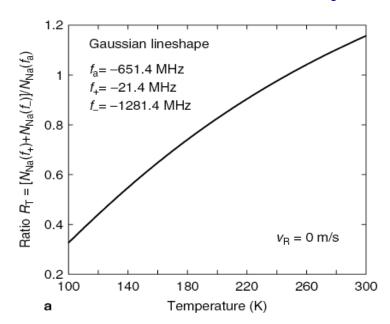


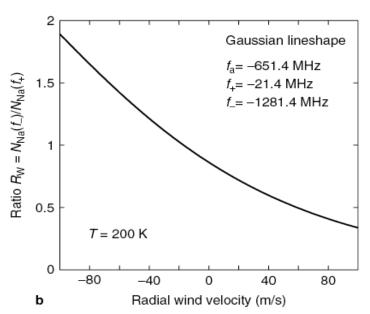
3-Frequency Doppler Ratio Technique

$$R_{T}(z) = \frac{N_{norm}(f_{+}, z, t_{1}) + N_{norm}(f_{-}, z, t_{2})}{N_{norm}(f_{a}, z, t_{3})}$$

$$\approx \frac{\sigma_{eff}(f_{+}, z) + \sigma_{eff}(f_{-}, z)}{\sigma_{eff}(f_{a}, z)}$$
(13.13)



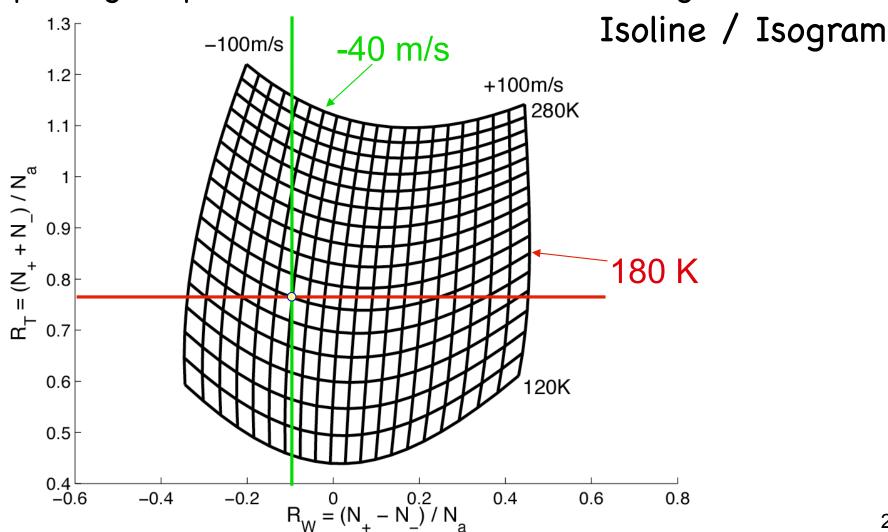






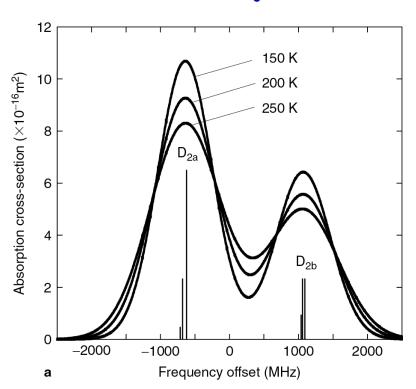
How Does Ratio Technique Work?

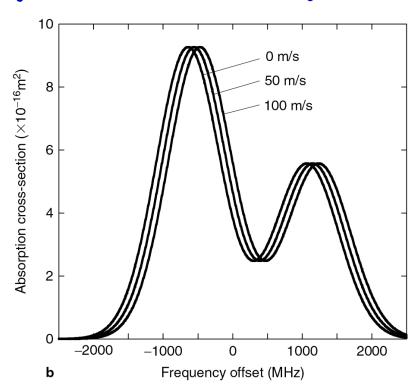
- Compute Doppler calibration curves from physics
- Look up these two ratios on the calibration curves to infer the corresponding Temperature and Wind from isoline/isogram.





Summary of Doppler Technique





Temperature determined from Doppler broadening width Radial wind determined from Doppler frequency shift

☐ Resonance absorption experiences 1-time Doppler shift and broadening, while Rayleigh scattering experiences 2 times of Doppler shift and broadening.



Summary

- ☐ The key point to measure temperature is to find and use temperature-dependent effects and phenomena to make measurements.
- Doppler technique utilizes the Doppler effect (frequency shift and linewidth broadening) by moving particles to infer wind and temperature information. It is widely applied in lidar, radar and sodar technique as well as passive optical remote sensing.
- Resonance fluorescence Doppler lidar technique applies scanning or ratio technique to infer the temperature and wind from the Doppler spectroscopy, while the Doppler spectroscopy is inferred from intensity ratio at different frequencies.