

Lecture 08. Fundamentals of Lidar Remote Sensing (4) "Solutions for Lidar Equations"

- Lidar Classification with Physical Processes
- Solution for scattering form lidar equation
- Solution for Raman lidar equation
- Solution for differential absorption lidar equation
- Solution for fluorescence form lidar equation
- Solution for resonance fluorescence lidar
- Solution for Rayleigh and Mie lidar in middle atmos

Summary



Review Lidar Equation

Lidar equation is to link the expected lidar returns (N_s) to the lidar parameters (both transmitter and receiver), transmission through medium, physical interactions between light and objects, and background/noise conditions, etc.

Keep in mind the big picture of a lidar system – Radiation source Radiation propagation in the medium Interaction of radiation with the objects Signal propagation in the medium Photons are collected, filtered and detected

Can you derive a lidar equation by yourself?



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Parameters in General Lidar Equation

Assumptions: independent and single scattering

$$N_{S}(\lambda,R) = \left[\frac{P_{L}(\lambda_{L})\Delta t}{hc/\lambda_{L}}\right] \cdot \left[\beta(\lambda,\lambda_{L},\theta,R)\Delta R\right] \cdot \frac{A}{R^{2}} \cdot \left[T(\lambda_{L},R)T(\lambda,R)\right] \cdot \left[\eta(\lambda,\lambda_{L})G(R)\right] + N_{B}$$

 N_s (R) – expected received photon number from a distance R (8.1)

- P_{L} transmitted laser power, λ_{L} laser wavelength
- Δt integration time,
- h Planck's constant, c light speed
- $\beta(R)$ volume scatter coefficient at distance R for angle θ ,
- ΔR thickness of the range bin
- A area of receiver,

T(R) – one way transmission of the light from laser source to distance R or from distance R to the receiver,

- η system optical efficiency,
- G(R) geometrical factor of the system,
- $N_{\rm B}$ background and detector noise photon counts.



General lidar equation with angular scattering coefficient

$$N_{S}(\lambda, R) = N_{L}(\lambda_{L}) \cdot \left[\beta(\lambda, \lambda_{L}, \theta, R)\Delta R\right] \cdot \frac{A}{R^{2}} \cdot \left[T(\lambda_{L}, R)T(\lambda, R)\right] \cdot \left[\eta(\lambda, \lambda_{L})G(R)\right] + N_{B}$$

General lidar equation with total scattering coefficient

$$N_{S}(\lambda,R) = N_{L}(\lambda_{L}) \cdot \left[\beta_{T}(\lambda,\lambda_{L},R)\Delta R\right] \cdot \frac{A}{4\pi R^{2}} \cdot \left[T(\lambda_{L},R)T(\lambda,R)\right] \cdot \left[\eta(\lambda,\lambda_{L})G(R)\right] + N_{B}$$

General lidar equation in angular scattering coefficient β and extinction coefficient α form

$$N_{S}(\lambda,R) = \left[\frac{P_{L}(\lambda_{L})\Delta t}{hc/\lambda_{L}}\right] \left[\beta(\lambda,\lambda_{L},\theta,R)\Delta R\right] \left(\frac{A}{R^{2}}\right)$$

$$\cdot \exp\left[-\int_{0}^{R} \alpha(\lambda_{L},r')dr'\right] \exp\left[-\int_{0}^{R} \alpha(\lambda,r')dr'\right] \left[\eta(\lambda,\lambda_{L})G(R)\right] + N_{B}$$

(4.2)

(4.12)



Specific Lidar Equations

Lidar equation for Rayleigh lidar

$$N_{S}(\lambda,R) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\beta(\lambda,R)\Delta R\right) \left(\frac{A}{R^{2}}\right) T^{2}(\lambda,R) \left(\eta(\lambda)G(R)\right) + N_{B}$$
(5.17)

Lidar equation for resonance fluorescence lidar

$$N_{S}(\lambda,R) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{eff}(\lambda,R)n_{c}(z)R_{B}(\lambda)\Delta R\right) \left(\frac{A}{4\pi R^{2}}\right) \left(T_{a}^{2}(\lambda,R)T_{c}^{2}(\lambda,R)\right) \left(\eta(\lambda)G(R)\right) + N_{B}(\lambda)G(R) + N_{B}(\lambda)G(R)$$

Lidar equation for differential absorption lidar

$$N_{S}(\lambda_{on}^{off}, R) = N_{L}(\lambda_{on}^{off}) \Big[\beta_{sca}(\lambda_{on}^{off}, R) \Delta R \Big] \Big(\frac{A}{R^{2}} \Big) \exp \Big[-2 \int_{0}^{z} \overline{\alpha}(\lambda_{on}^{off}, r') dr' \Big] \\ \times \exp \Big[-2 \int_{0}^{z} \sigma_{abs}(\lambda_{on}^{off}, r') n_{c}(r') dr' \Big] \Big[\eta(\lambda_{on}^{off}) G(R) \Big] + N_{B}$$
(5.21-5.22)



Elastic Scattering Lidar: Rayleigh Lidar and Mie Lidar



Block diagram of four-wavelength lidar.



Solution for Scattering Form Lidar Equation

□ Scattering form lidar equation

$$N_{S}(\lambda,R) = \left[\frac{P_{L}(\lambda_{L})\Delta t}{hc/\lambda_{L}}\right] \cdot \left[\beta(\lambda,\lambda_{L},R)\Delta R\right] \cdot \left[\frac{A}{R^{2}}\right] \cdot \left[T(\lambda_{L},R)T(\lambda,R)\right] \cdot \left[\eta(\lambda,\lambda_{L})G(R)\right] + N_{B}$$

□ Solution for scattering form lidar equation

$$\beta(\lambda,\lambda_L,R) = \frac{N_S(\lambda,R) - N_B}{\left[\frac{P_L(\lambda_L)\Delta t}{hc/\lambda_L}\right] \Delta R \left(\frac{A}{R^2}\right) \left[T(\lambda_L,R)T(\lambda,R)\right] \left[\eta(\lambda,\lambda_L)G(R)\right]}$$
(8.3)

(8.2)



Inelastic Scattering Lidar: Raman Lidar

Raman spectra for O₂, N₂ and H₂O under 355nm laser wavelength





H₂O in atmosphere

Fig. 9.3. Typical setup of a water-vapor and aerosol Raman lidar. AL – achromatic lens, IF – interference filter, DBS – dichroic beam splitter.

Fig. 9.2. Raman backscatter spectrum of the atmosphere for an incident laser wavelength of 355 nm, normal pressure, a temperature of 300 K, an N₂ and O₂ content of 0.781 and 0.209, respectively, and a water-vapor mixing ratio of 10 g/kg. The curves for liquid water and ice are arbitrarily scaled. The isosbestic point is discussed in Subsection 9.5.2.



Solution for Raman Lidar Equation

□ Raman lidar equations for gas of interest (8.4) and reference gas (8.5)

$$P_{cRa}(R,\lambda_{cRa}) = \frac{E_o \eta_{\lambda_{cRa}}}{R^2} O(R,\lambda_{cRa}) \beta_{cRa}(R,\lambda_0,\lambda_{cRa}) \exp\left\{-\int_0^R \left[\alpha(r,\lambda_0) + \alpha(r,\lambda_{cRa})\right] dr\right\} (8.4)$$

$$P_{\text{Ref}Ra}(R,\lambda_{\text{Ref}Ra}) = \frac{E_o \eta_{\lambda_{\text{Ref}Ra}}}{R^2} O(R,\lambda_{\text{Ref}Ra}) \beta_{\text{Ref}Ra}(R,\lambda_0,\lambda_{\text{Ref}Ra}) \exp\left\{-\int_0^R \left[\alpha(r,\lambda_0) + \alpha(r,\lambda_{\text{Ref}Ra})\right] dr\right\}$$
(8.5)

Solution for Raman lidar equations

$$\frac{P_{cRa}(R,\lambda_{cRa})}{P_{\text{Ref}Ra}(R,\lambda_{\text{Ref}Ra})} = \frac{\eta_{\lambda_{cRa}}}{\eta_{\lambda_{\text{Ref}Ra}}} \frac{N_{cRa}(R) \frac{d\sigma_{cRa}}{d\Omega}(\pi,\lambda_0,\lambda_{cRa})}{N_{\text{Ref}Ra}(R) \frac{d\sigma_{\text{Ref}Ra}}{d\Omega}(\pi,\lambda_0,\lambda_{\text{Ref}Ra})} \frac{\exp\left\{-\int_0^R \alpha(r,\lambda_{cRa})dr\right\}}{\exp\left\{-\int_0^R \alpha(r,\lambda_{\text{Ref}Ra})dr\right\}}$$

$$Mixing Ratio = \frac{N_{cRa}(R)}{N_{\text{Ref}Ra}(R)} \tag{8.6}$$

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Fig. 4.10. The differential absorption lidar technique

Collis and Russell in "Laser monitoring of the atmosphere"



Differential Absorption/Scattering Form

 $\hfill \label{eq:linear}$ For the laser with wavelength λ_{on} on the molecular absorption line

$$N_{S}(\lambda_{on},R) = N_{L}(\lambda_{on}) \left[\beta_{sca}(\lambda_{on},R)\Delta R\right] \left(\frac{A}{R^{2}}\right) \exp\left[-2\int_{0}^{R}\overline{\alpha}(\lambda_{on},r')dr'\right]$$
$$\times \exp\left[-2\int_{0}^{R}\sigma_{abs}(\lambda_{on},r')n_{c}(r')dr'\right] \left[\eta(\lambda_{on})G(R)\right] + N_{B}$$

 $\hfill \label{eq:linear}$ For the laser with wavelength λ_{off} off the molecular absorption line

$$N_{S}(\lambda_{off},R) = N_{L}(\lambda_{off}) \Big[\beta_{sca}(\lambda_{off},R) \Delta R \Big] \Big(\frac{A}{R^{2}} \Big) \exp \Big[-2 \int_{0}^{R} \overline{\alpha}(\lambda_{off},r') dr' \Big] \\ \times \exp \Big[-2 \int_{0}^{R} \sigma_{abs}(\lambda_{off},r') n_{c}(r') dr' \Big] \Big[\eta(\lambda_{off}) G(R) \Big] + N_{B}$$

$$(8.8)$$

(8.7)



Differential Absorption/Scattering Form

□ The ratio of photon counts from these two channels is a function of the differential absorption and scattering:

$$\frac{N_{S}(\lambda_{on},R) - N_{B}}{N_{S}(\lambda_{off},R) - N_{B}} = \frac{N_{L}(\lambda_{on})\beta_{sca}(\lambda_{on},R)}{N_{L}(\lambda_{off})\beta_{sca}(\lambda_{off},R)} \frac{\eta(\lambda_{on})}{\eta(\lambda_{off})} \times \exp\left\{-2\int_{0}^{R} \left[\overline{\alpha}(\lambda_{on},r') - \overline{\alpha}(\lambda_{off},r')\right]dr'\right\} \times \exp\left\{-2\int_{0}^{R} \left[\sigma_{abs}(\lambda_{on},r') - \sigma_{abs}(\lambda_{off},r')\right]n_{c}(r')dr'\right\}$$
(8.9)

$$\Delta \sigma = \sigma_{abs}(\lambda_{on}) - \sigma_{abs}(\lambda_{off})$$
(8.10)



Solution for Differential Absorption Lidar Equation

Solution for differential absorption lidar equation

$$n_{c}(R) = \frac{1}{2\Delta\sigma} \frac{d}{dR} \begin{cases} \ln \left[\frac{N_{L}(\lambda_{on})\beta_{sca}(\lambda_{on},R)}{N_{L}(\lambda_{off})\beta_{sca}(\lambda_{off},R)} \frac{\eta(\lambda_{on})}{\eta(\lambda_{off})} \right] \\ -\ln \left[\frac{N_{S}(\lambda_{on},R) - N_{B}}{N_{S}(\lambda_{off},R) - N_{B}} \right] \\ -2\int_{0}^{R} \left[\overline{\alpha}(\lambda_{on},r') - \overline{\alpha}(\lambda_{off},r') \right] dr' \end{cases}$$
(8.11)

$$\Delta \sigma = \sigma_{abs}(\lambda_{on}) - \sigma_{abs}(\lambda_{off}) \qquad (8.10)$$



Resonance Fluorescence Lidar













Solution for Fluorescence Form Lidar Equation

□ Fluorescence form lidar equation

$$N_{S}(\lambda,R) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{eff}(\lambda,R)n_{c}(R)R_{B}(\lambda)\Delta R\right) \left(\frac{A}{4\pi R^{2}}\right) \left(T_{a}^{2}(\lambda,R)T_{c}^{2}(\lambda,R)\right) \left(\eta(\lambda)G(R)\right) + N_{B}(\lambda)G(R) + N_{B}(\lambda)G(R)$$

□ Solution for fluorescence form lidar equation

$$\begin{split} n_{c}(R) &= \frac{N_{S}(\lambda,R) - N_{B}}{\left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{eff}(\lambda)R_{B}(\lambda)\Delta R\right) \left(\frac{A}{4\pi R^{2}}\right) \left(\eta(\lambda)T_{a}^{2}(\lambda,R)T_{c}^{2}(\lambda,R)G(R)\right)} \end{split}$$

(8.13)



Solution for

Resonance Fluorescence Lidar Equation

Resonance fluorescence and Rayleigh lidar equations

 $N_{S}(\lambda,z) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{eff}(\lambda,z)n_{c}(z)R_{B}(\lambda)\Delta z\right) \left(\frac{A}{4\pi z^{2}}\right) \left(T_{a}^{2}(\lambda)T_{c}^{2}(\lambda,z)\right) \left(\eta(\lambda)G(z)\right) + N_{B}$ (8.12) $N_{R}(\lambda,z_{R}) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\sigma_{R}(\pi,\lambda)n_{R}(z_{R})\Delta z\right) \left(\frac{A}{z_{R}^{2}}\right) T_{a}^{2}(\lambda,z_{R}) \left(\eta(\lambda)G(z_{R})\right) + N_{B}$ (8.14)

Rayleigh normalization

 $\frac{n_c(z)}{n_R(z_R)} = \frac{N_S(\lambda, z) - N_B}{N_R(\lambda, z_R) - N_B} \cdot \frac{z^2}{z_R^2} \cdot \frac{4\pi\sigma_R(\pi, \lambda)}{\sigma_{eff}(\lambda, z)R_B(\lambda)} \cdot \frac{T_a^2(\lambda, z_R)G(z_R)}{T_a^2(\lambda, z)T_c^2(\lambda, z)G(z)}$ (8.15)

Solution for resonance fluorescence

$$n_{c}(z) = n_{R}(z_{R}) \frac{N_{S}(\lambda, z) - N_{B}}{N_{R}(\lambda, z_{R}) - N_{B}} \cdot \frac{z^{2}}{z_{R}^{2}} \cdot \frac{4\pi\sigma_{R}(\pi, \lambda)}{\sigma_{eff}(\lambda, z)R_{B}(\lambda)} \cdot \frac{1}{T_{c}^{2}(\lambda, z)}$$
(8.16)



Rayleigh and Mie Lidar for Middle Atmosphere Detection



Polar mesospheric clouds (PMC), also noctilucent clouds (NLC), are thin scattering layers of nanometer-sized water ice particles, occurring around 80-87 km in the high-latitude summer mesopause region.





Solution for Rayleigh and Mie Lidars

Rayleigh and Mie (middle atmos) lidar equations

$$N_{S}(\lambda,z) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\beta_{R}(z) + \beta_{aerosol}(z)\right) \Delta z \left(\frac{A}{z^{2}}\right) T_{a}^{2}(\lambda,z) \left(\eta(\lambda)G(z)\right) + N_{B} \quad (8.17)$$

$$N_{R}(\lambda,z_{R}) = \left(\frac{P_{L}(\lambda)\Delta t}{hc/\lambda}\right) \left(\beta_{R}(z_{R})\Delta z\right) \left(\frac{A}{z_{R}^{2}}\right) T_{a}^{2}(\lambda,z_{R}) \left(\eta(\lambda)G(z_{R})\right) + N_{B} \quad (8.18)$$

Rayleigh normalization

$$\frac{\beta_R(z) + \beta_{aerosol}(z)}{\beta_R(z_R)} = \frac{N_S(\lambda, z) - N_B}{N_R(\lambda, z_R) - N_B} \cdot \frac{z^2}{z_R^2} \cdot \frac{T_a^2(\lambda, z_R)G(z_R)}{T_a^2(\lambda, z)G(z)}$$
(8.19)

\Box For Rayleigh scattering at z and z_R

$$\frac{\beta_R(z)}{\beta_R(z_R)} = \frac{\sigma_R(z)n_{atm}(z)}{\sigma_R(z_R)n_{atm}(z_R)} = \frac{n_{atm}(z)}{n_{atm}(z_R)}$$
(8.20)

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Solution (Continued)

Solution for Mie scattering in middle atmosphere

$$\beta_{aerosol}(z) = \beta_R(z_R) \left[\frac{N_S(\lambda, z) - N_B}{N_R(\lambda, z_R) - N_B} \cdot \frac{z^2}{z_R^2} - \frac{n_{atm}(z)}{n_{atm}(z_R)} \right]$$
(8.21)
$$\beta_R(\lambda, z_R, \pi) = 2.938 \times 10^{-32} \frac{P(z_R)}{T(z_R)} \cdot \frac{1}{\lambda^{4.0117}} \left(m^{-1} s r^{-1} \right)$$
(5.14)

$$\frac{\beta_R(z)}{\beta_R(z_R)} = \frac{N_S(\lambda, z) - N_B}{N_R(\lambda, z_R) - N_B} \cdot \frac{z^2}{z_R^2} \cdot \frac{T_a^2(\lambda, z_R)G(z_R)}{T_a^2(\lambda, z)G(z_R)}$$
(8.22)

Solution for relative number density in Rayleigh lidar

$$RND(z) = \frac{n_{atm}(z)}{n_{atm}(z_R)} = \frac{\beta_R(z)}{\beta_R(z_R)} = \frac{N_S(\lambda, z) - N_B}{N_R(\lambda, z_R) - N_B} \cdot \frac{z^2}{z_R^2}$$
(8.23)

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Rayleigh Backscatter Cross Section

□ It is common in lidar field to calculate the Rayleigh backscatter cross section using the following equation

$$\frac{d\sigma_m(\lambda)}{d\Omega} = 5.45 \cdot \left(\frac{550}{\lambda}\right)^4 \times 10^{-32} \left(m^2 s r^{-1}\right)$$
(5.16)

where λ is the wavelength in nm.

□ The Rayleigh backscatter cross section can also be estimated from the Rayleigh backscatter coefficient

$$\beta_{Rayleigh}(\lambda, z, \theta = \pi) = 2.938 \times 10^{-32} \frac{P(z)}{T(z)} \cdot \frac{1}{\lambda^{4.0117}} \left(m^{-1} s r^{-1} \right)$$
(5.14)

where λ is the wavelength in meter, P in mbar, T in Kelvin.

$$\therefore \frac{d\sigma_m(\lambda)}{d\Omega} = \frac{\beta_{Rayleigh}(\lambda, z, \pi)}{n_{atmos}(z)} \left(m^2 s r^{-1}\right)$$
(8.24)



Summary

Physical processes in the interaction between radiation and objects classify lidars into different kinds, each possessing unique features and lidar equations along with certain applications.

Solutions of lidar equations can be obtained by solving the lidar equations directly if all the lidar parameters and atmosphere conditions are well known.

Solutions for three forms of lidar equations are shown: scattering form, fluorescence form, and differential absorption form, along with a few specific lidar types.

However, system parameters and atmosphere conditions may vary frequently and are NOT well known to experimenters.

□ A good solution is to perform Rayleigh normalization to cancel out most of the system and atmosphere parameters so that the essential and known parts can be solved.