



Lecture 11. Temperature Lidar (1)

Overview and Doppler Technique

- ☐ Overview of Temperature Measurement Techniques
- ☐ Doppler Technique for Temperature and Wind Measurements
- ☐ Resonance Fluorescence Na Doppler Lidar
- ☐ Summary



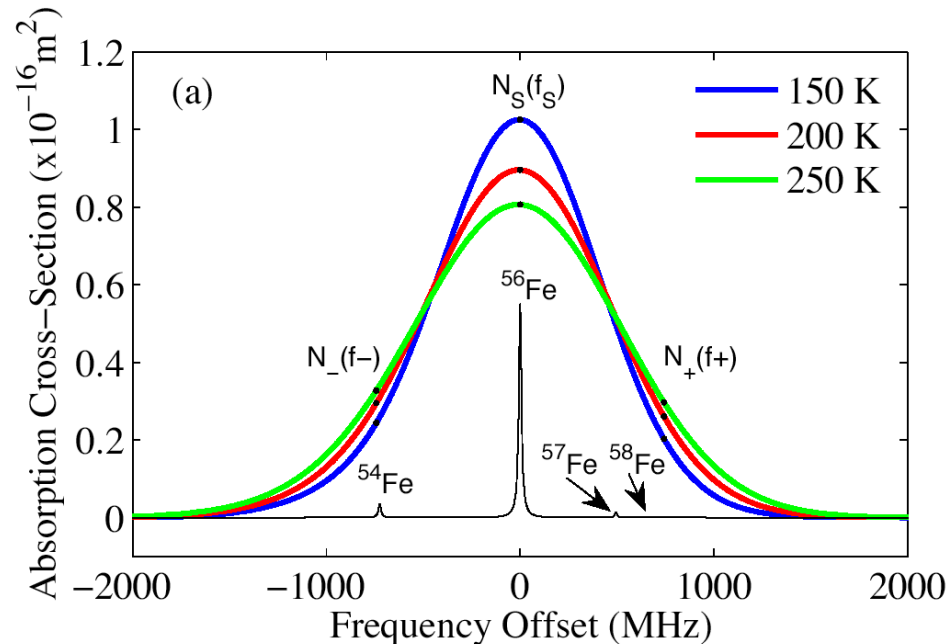
Temperature Measurement Techniques

Use temperature-dependent effects or phenomena

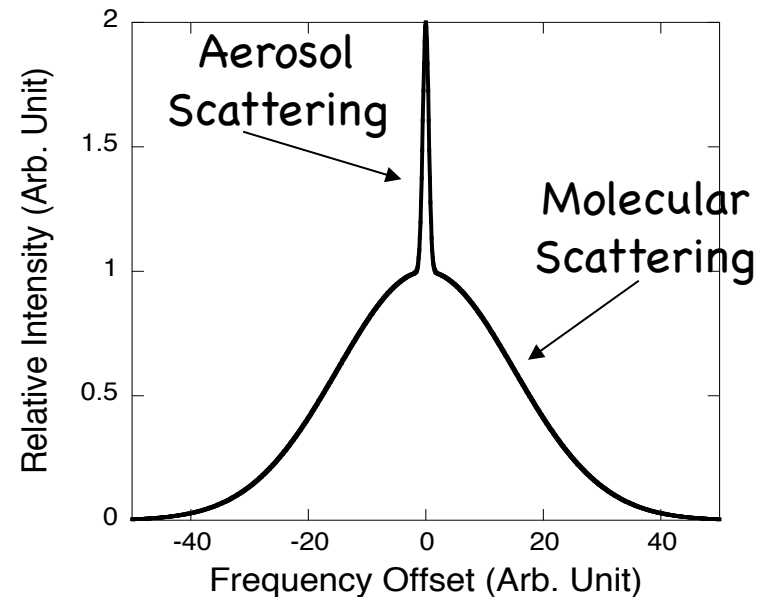
- ❑ **Doppler Technique** - Doppler broadening (not only for Na, K, and Fe, but also for Rayleigh scattering, as long as Doppler broadening dominate and can be detected)
- ❑ **Boltzmann Technique** - Boltzmann distribution of atomic populations on different energy levels (not only for Fe, but also for molecular spectroscopy in optical remote sensing)
- ❑ **Integration Technique (Rayleigh or Raman)** - integration lidar technique using ideal gas law and assuming hydrostatic equilibrium (not only for modern lidar, but also for cw searchlight and rocket falling sphere - some way to measure atmosphere number density)
- ❑ **Rotational Raman Technique** - temperature dependence of population ratio, similar to Boltzmann technique



Overview: Doppler Technique



Atomic Absorption Line



Rayleigh Scattering

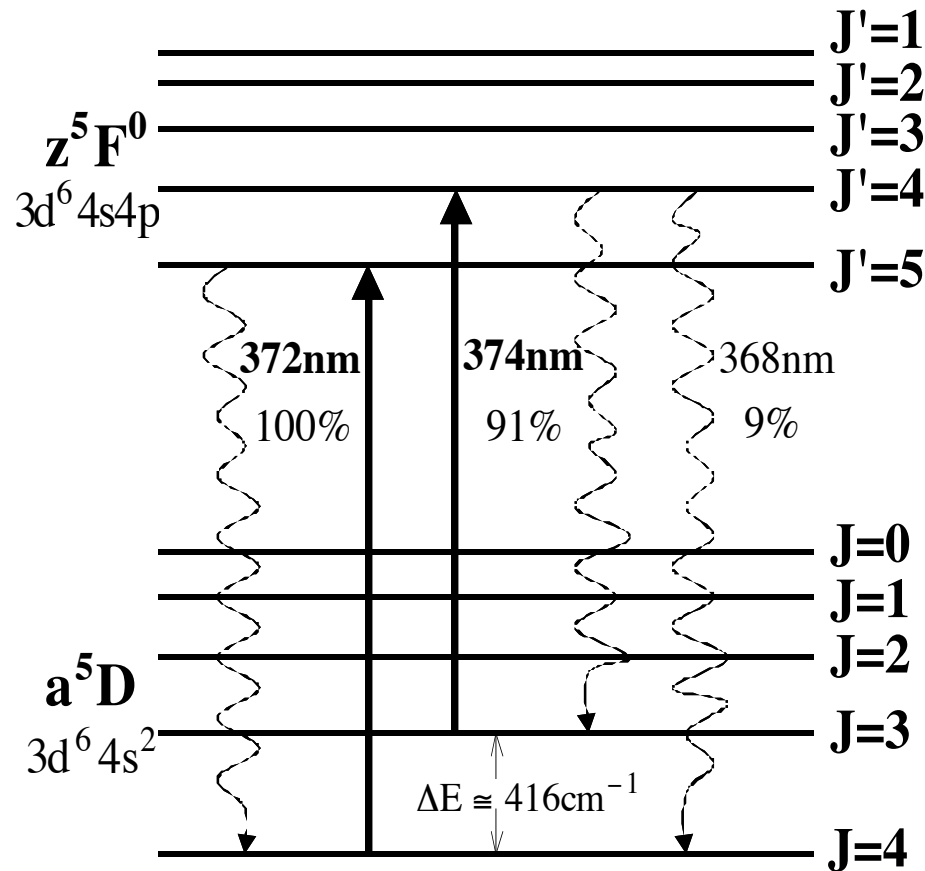
$$\sigma_{rms} = \frac{\nu_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}} = \sqrt{\frac{k_B T}{M \lambda_0^2}}$$

$$\sigma_{rms} = 2\nu_0 / c \sqrt{k_B T / M} = \frac{2}{\lambda_0} \sqrt{k_B T / M}$$

Doppler Spectrum (Doppler Broadening Width) \Rightarrow Temperature



Overview: Boltzmann Technique



Atomic Fe Energy Level

[Gelbwachs, 1994; Chu et al., 2002]

**Maxwell-Boltzmann Distribution
in Thermal-dynamic Equilibrium**

$$\frac{P_2(J=3)}{P_1(J=4)} = \frac{\rho_{Fe(374)}}{\rho_{Fe(372)}} = \frac{g_2}{g_1} \exp(-\Delta E / k_B T)$$



$$T = \frac{\Delta E / k_B}{\ln \left(\frac{g_2 \cdot P_1}{g_1 \cdot P_2} \right)}$$

P_1, P_2 -- Fe populations

g_1, g_2 -- Degeneracy

k_B -- Boltzmann constant

T -- Temperature

Population Ratio \Rightarrow Temperature



Overview: Integration Technique

Hydrostatic Equation

$$dP = -\rho g dz$$

+

Ideal Gas Law

$$P = \rho R T$$

$$T(z) = T(z_o) \frac{\rho(z_o)}{\rho(z)} + \frac{1}{R} \int_z^{z_o} g(r) \frac{\rho(r)}{\rho(z)} dr$$

Seeding
Temperature

Relative
Density

$T(z_0)$ - Seeding Temperature;

R - gas constant for dry air;

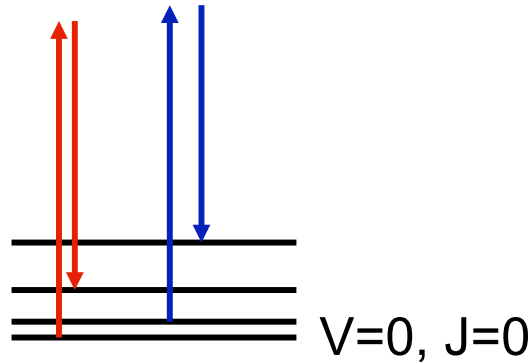
ρ - number density

g - gravitational acceleration

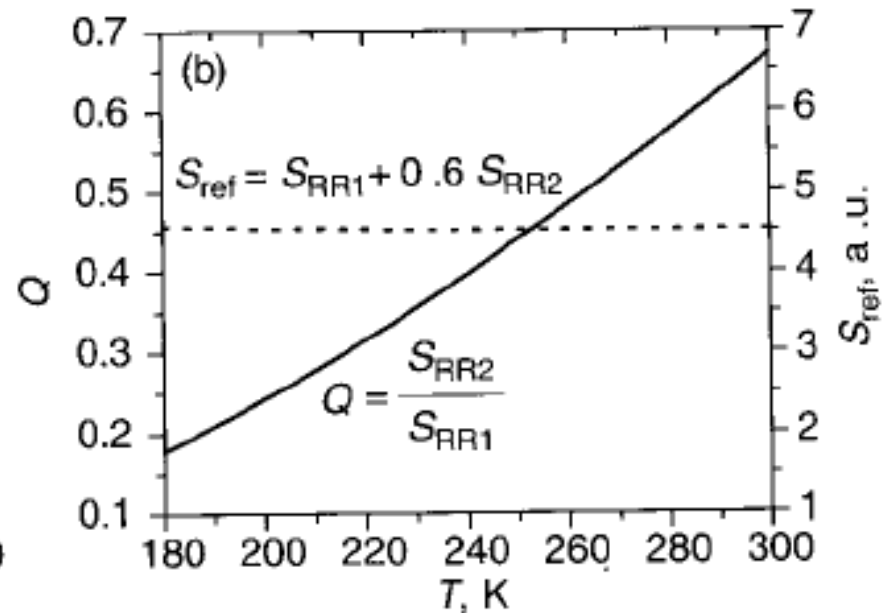
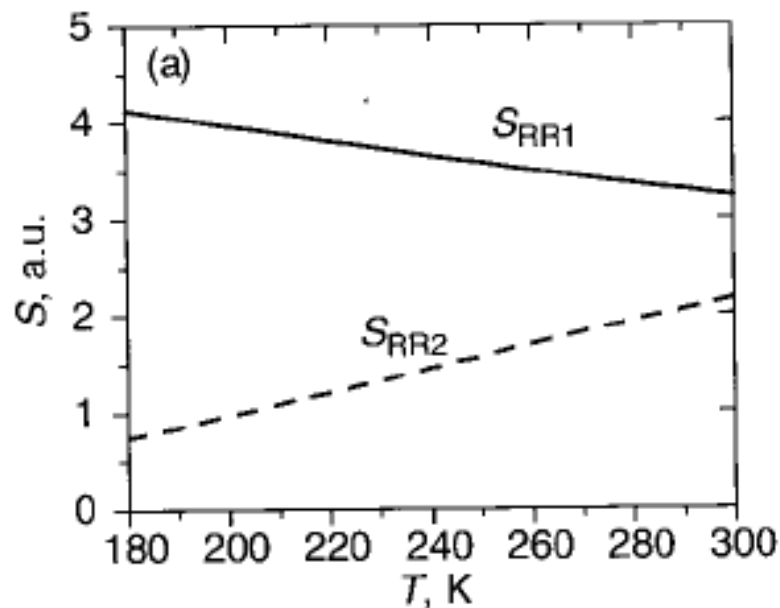
**Number Density Ratio \Rightarrow Temperature
(lidar backscatter ratio at different altitudes)**



Overview: Rotation Raman Technique



$$Q(T) = \frac{\sum_{i=O_2, N_2} \sum_{J_i} \tau_{RR2}(J_i) \eta_i \left(\frac{d\sigma}{d\Omega} \right)_{\pi}^{RR,i}(J_i)}{\sum_{i=O_2, N_2} \sum_{J_i} \tau_{RR1}(J_i) \eta_i \left(\frac{d\sigma}{d\Omega} \right)_{\pi}^{RR,i}(J_i)}$$



□ Temperature can be derived from the ratio of two pure Rotational Raman line intensities. This is essentially the same principle as Boltzmann temperature technique!

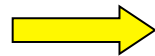


Doppler Technique to Measure Temperature and Wind

□ Doppler effect is commonly experienced by moving particles, such as atoms, molecules, and aerosols. It is the apparent frequency change of radiation that is perceived by the particles moving relative to the source of the radiation. This is called Doppler shift.

□ Doppler frequency shift is proportional to the radial velocity along the line of sight (LOS) of the radiation -

$$\omega = \omega_0 - \vec{k} \cdot \vec{v}$$



$$\begin{aligned}\Delta\omega &= \omega - \omega_0 = -\vec{k} \cdot \vec{v} = -\omega_0(v/c)\cos\theta \\ \Delta\nu &= -\nu_0(v/c)\cos\theta = -(\nu/\lambda_0)\cos\theta\end{aligned}$$

where ω_0 is the radiation frequency at rest, ω is the shifted frequency, k is the wave vector of the radiation ($k=2\pi/\lambda$), and v is the particle velocity.



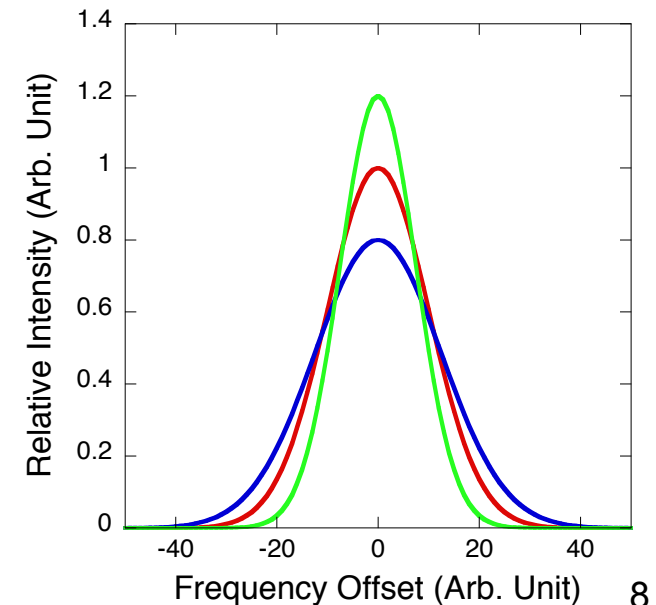
Doppler Technique to Measure Temperature and Wind

□ Due to particles' thermal motions in the atmosphere, the distribution of perceived frequencies for all particles mirrors their velocity distribution. According to the Maxwellian velocity distribution, the perceived frequencies by moving particles has a Gaussian lineshape, given by

$$\exp\left(-\frac{Mv_z^2}{2k_B T}\right)dv_z = \exp\left\{-\frac{Mc^2(v-v_0)^2}{2v_0^2 k_B T}\right\} \frac{c}{v_0} dv$$

□ The peak is at $\omega = \omega_0$ and the rms width is give by

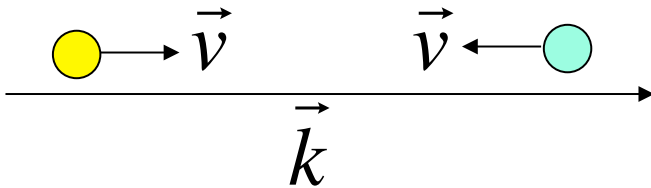
$$\sigma_{rms} = \frac{v_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}}$$





Doppler Shift in Resonance Absorption

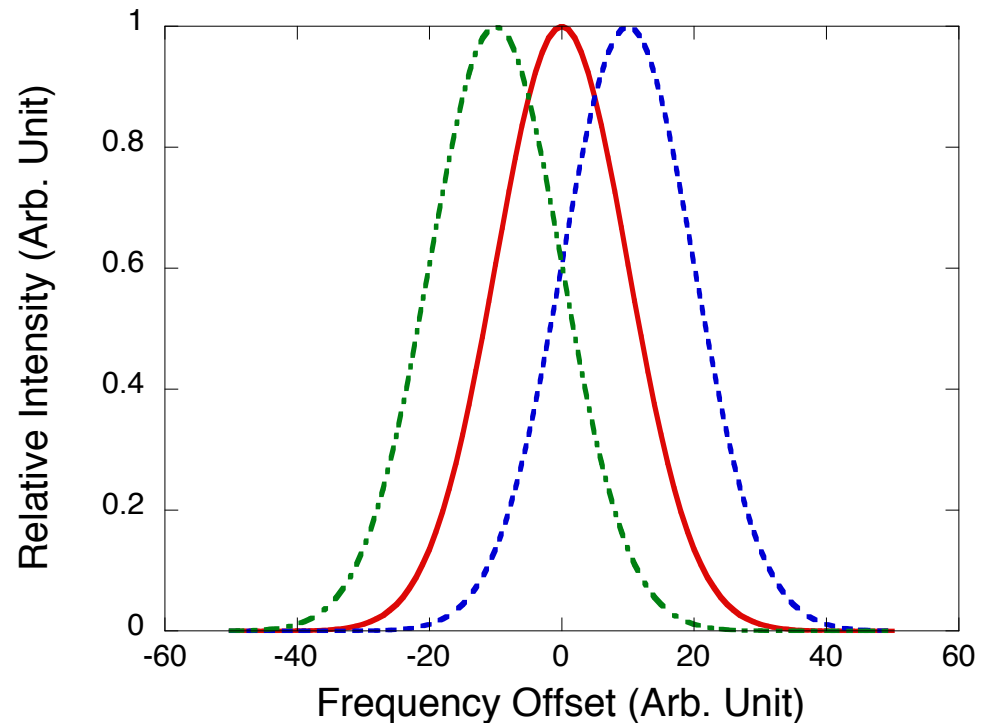
$$\Delta\omega = \omega - \omega_0 = -\vec{k} \cdot \vec{v} = -\omega_0 \frac{v \cos \theta}{c}$$



Emitter and receiver move towards each other:

-Blue shift in perceived radiation frequency

-Red shift in absorption peak frequency



□ The velocity measurements of lidar, radar, and sodar all base on the Doppler shift principle !

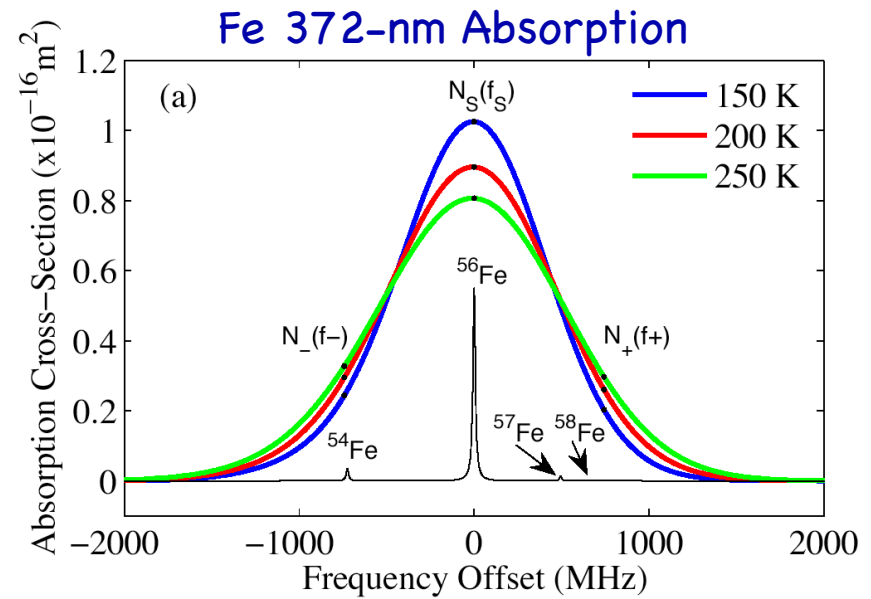
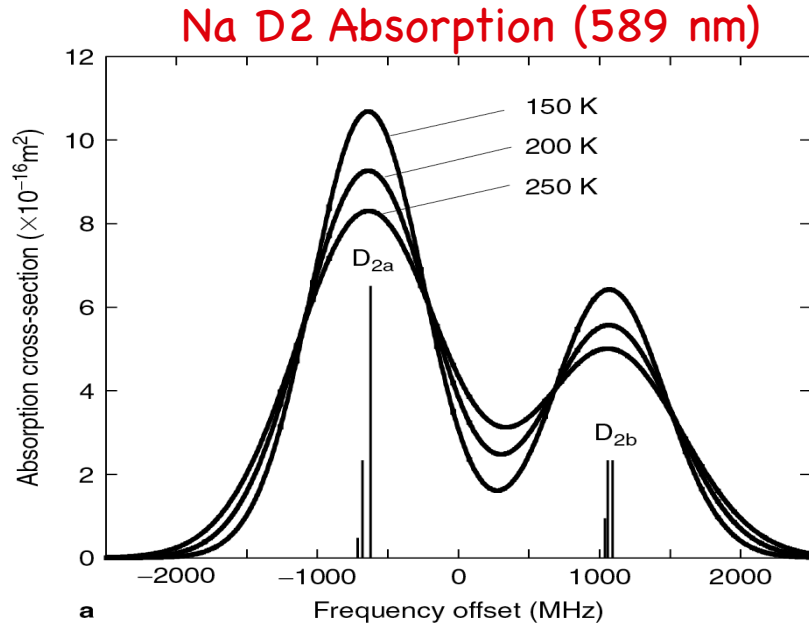


Doppler Broadening in Resonance Absorption Lines

$$\sigma_{rms} = \frac{\nu_0}{c} \sqrt{\frac{k_B T}{M}} = \frac{1}{\lambda_0} \sqrt{\frac{k_B T}{M}}$$

$T \uparrow \Rightarrow \sigma_{rms} \uparrow$

$M \uparrow \Rightarrow \sigma_{rms} \downarrow$





Doppler Shift in Rayleigh Scattering

- Refer to textbook 5.2.2.4 Lidar wind vs radar wind measurements

Momentum Conservation $m\vec{v}_1 + \hbar\vec{k}_1 = m\vec{v}_2 + \hbar\vec{k}_2$

Energy Conservation $\frac{1}{2}mv_1^2 + \hbar\omega_1 = \frac{1}{2}mv_2^2 + \hbar\omega_2$

}

$$\omega_1 = \omega_2 + \vec{k}_1 \cdot \vec{v}_1 - \vec{k}_2 \cdot \vec{v}_2 + \frac{\hbar k_1^2}{2m} - \frac{\hbar k_2^2}{2m}$$

- For Rayleigh or radar backscatter signals, we have

$$\vec{k}_2 \approx -\vec{k}_1 \quad \vec{v}_2 \approx \vec{v}_1$$

- The frequency shift for Rayleigh or radar backscattering is

$$\Delta\omega_{\text{Rayleigh, backscatter}} = \omega_2 - \omega_1 = -2\vec{k}_1 \cdot \vec{v}_1$$



Doppler Broadening in Rayleigh Scatter

□ To derive the Doppler broadening, let's write the Doppler shift as

$$\omega = \omega_0 \left(1 - \frac{2v_R}{c} \right) \quad \longrightarrow \quad v_R = \frac{\omega_0 - \omega}{2\omega_0 / c} = \frac{\nu_0 - \nu}{2\nu_0 / c}$$

□ According to the Maxwellian velocity distribution, the relative probability that an atom/molecule in a gas at temperature T has its velocity component along the line of sight between v_R and $v_R + dv_R$ is

$$P(v_R \rightarrow v_R + dv_R) \propto \exp(-Mv_R^2 / 2k_B T) dv_R$$

□ Substitute the v_R equation into the Maxwellian distribution,

$$I \propto \exp\left(-\frac{M(\nu_0 - \nu)^2}{2k_B T (2\nu_0 / c)^2}\right) (c / 2\nu_0) d\nu$$

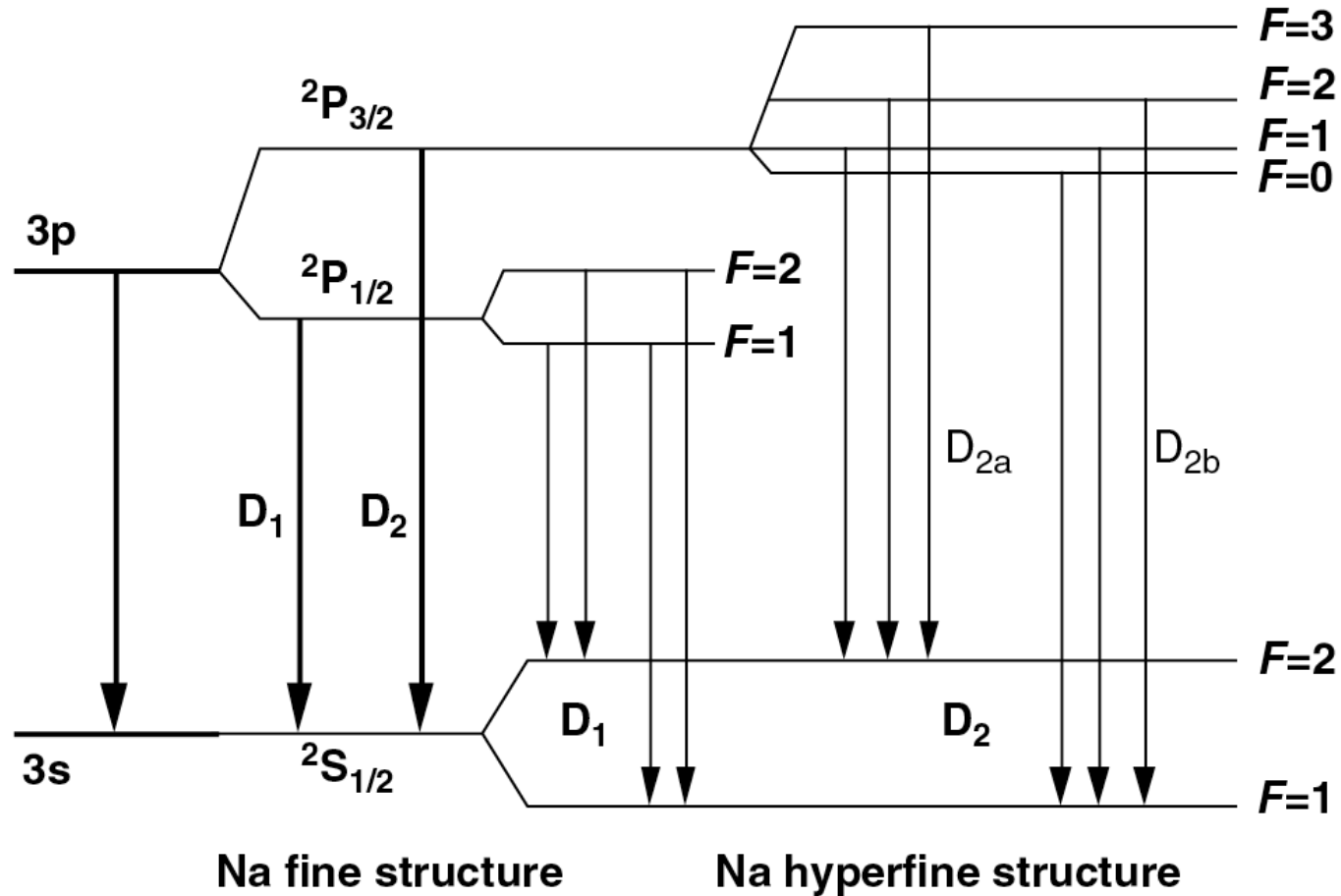
□ Therefore, the rms width of the Doppler broadening is

$$\sigma_{rms} = 2\nu_0 / c \sqrt{k_B T / M} = \frac{2}{\lambda_0} \sqrt{k_B T / M}$$

2 times !

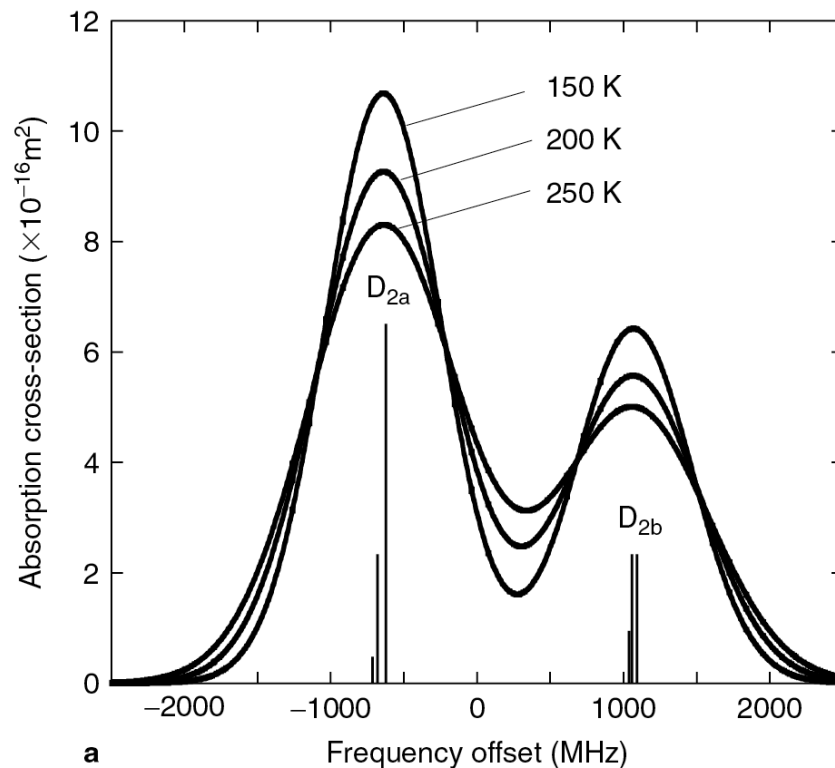


Na Atomic Energy Levels

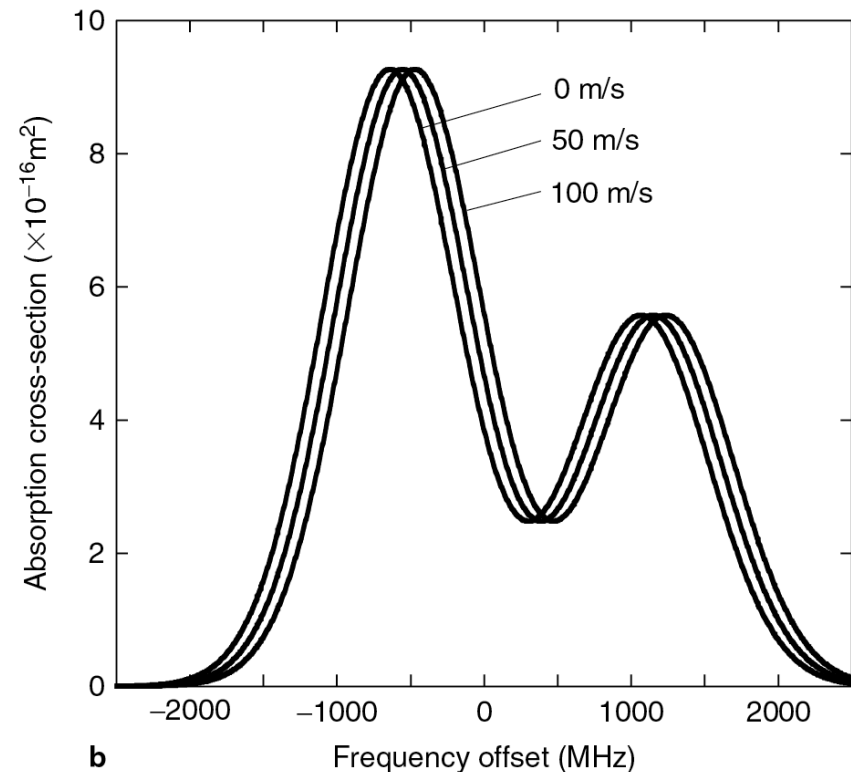




Doppler Effect in Na D₂ Line Resonance Fluorescence



Na D₂ absorption linewidth is temperature dependent



Na D₂ absorption peak freq is wind dependent



Na Atomic Parameters

Table 5.1 Parameters of the Na D₁ and D₂ Transition Lines

Transition Line	Central Wavelength (nm)	Transition Probability (10 ⁸ s ⁻¹)	Radiative Lifetime (nsec)	Oscillator Strength f_{ik}
D ₁ (² P _{1/2} → ² S _{1/2})	589.7558	0.614	16.29	0.320
D ₂ (² P _{3/2} → ² S _{1/2})	589.1583	0.616	16.23	0.641
Group	² S _{1/2}	² P _{3/2}	Offset (GHz)	Relative Line Strength ^a
D _{2b}	$F = 1$	$F = 2$	1.0911	5/32
		$F = 1$	1.0566	5/32
		$F = 0$	1.0408	2/32
D _{2a}	$F = 2$	$F = 3$	−0.6216	14/32
		$F = 2$	−0.6806	5/32
		$F = 1$	−0.7150	1/32
Doppler-Free Saturation–Absorption Features of the Na D ₂ Line				
f_a (MHz)	f_c (MHz)	f_b (MHz)	f_+ (MHz)	f_- (MHz)
−651.4	187.8	1067.8	−21.4	−1281.4

^aRelative line strengths are in the absence of a magnetic field or the spatial average. When Hanle effect is considered in the atmosphere, the relative line strengths will be modified depending on the geomagnetic field and the laser polarization.



Doppler-Limited Na Spectroscopy

- Doppler-broadened Na absorption cross-section is approximated as a Gaussian with rms width σ_D

$$\sigma_{abs}(\nu) = \frac{1}{\sqrt{2\pi}\sigma_D} \frac{e^2 f}{4\epsilon_0 m_e c} \sum_{n=1}^6 A_n \exp\left(-\frac{[\nu_n - \nu(1 - V_R/c)]^2}{2\sigma_D^2}\right)$$

- Assume the laser lineshape is a Gaussian with rms width σ_L
- The effective cross-section is the convolution of the atomic absorption cross-section and the laser lineshape

$$\sigma_{eff}(\nu) = \frac{1}{\sqrt{2\pi}\sigma_e} \frac{e^2 f}{4\epsilon_0 m_e c} \sum_{n=1}^6 A_n \exp\left(-\frac{[\nu_n - \nu(1 - V_R/c)]^2}{2\sigma_e^2}\right)$$

where $\sigma_e = \sqrt{\sigma_D^2 + \sigma_L^2}$ and $\sigma_D = \sqrt{\frac{k_B T}{M \lambda_0^2}}$

The frequency discriminator/analyzer is in the atmosphere!!! 16



Doppler Scanning Technique

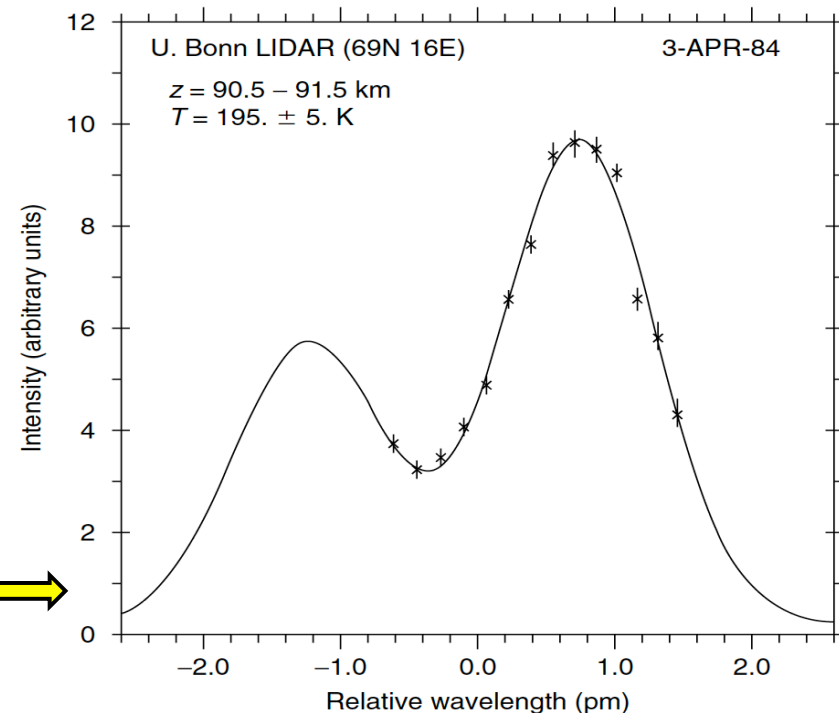
$$N_{Na}(\lambda, z) = \left(\frac{P_L(\lambda) \Delta t}{hc/\lambda} \right) (\sigma_{eff}(\lambda) n_{Na}(z) \Delta z) \left(\frac{A}{4\pi z^2} \right) (\eta(\lambda) T_a^2(\lambda) T_c^2(\lambda, z) G(z))$$

$$N_R(\lambda, z_R) = \left(\frac{P_L(\lambda) \Delta t}{hc/\lambda} \right) (\sigma_R(\pi, \lambda) n_R(z_R) \Delta z) \left(\frac{A}{z_R^2} \right) (\eta(\lambda) T_a^2(\lambda, z_R) G(z_R))$$

$$\sigma_{eff}(\lambda, z) = \frac{C(z)}{T_c^2(\lambda, z)} \frac{N_{Na}(\lambda, z)}{N_R(\lambda, z_R)}$$

where $C(z) = \frac{\sigma_R(\pi, \lambda) n_R(z_R)}{n_{Na}(z)} \frac{4\pi z^2}{z_R^2}$

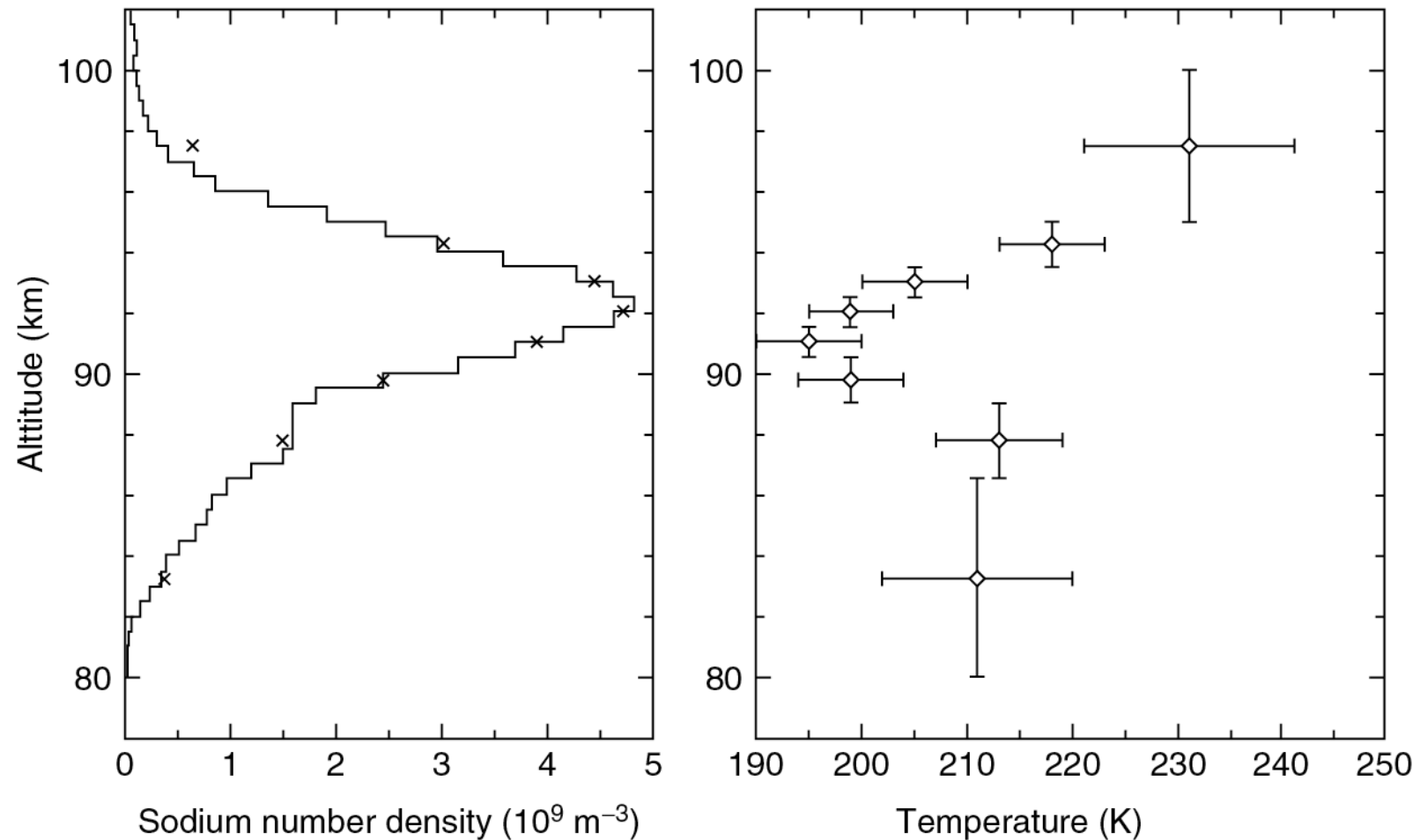
[Fricke and von Zahn, JATP, 1985] →





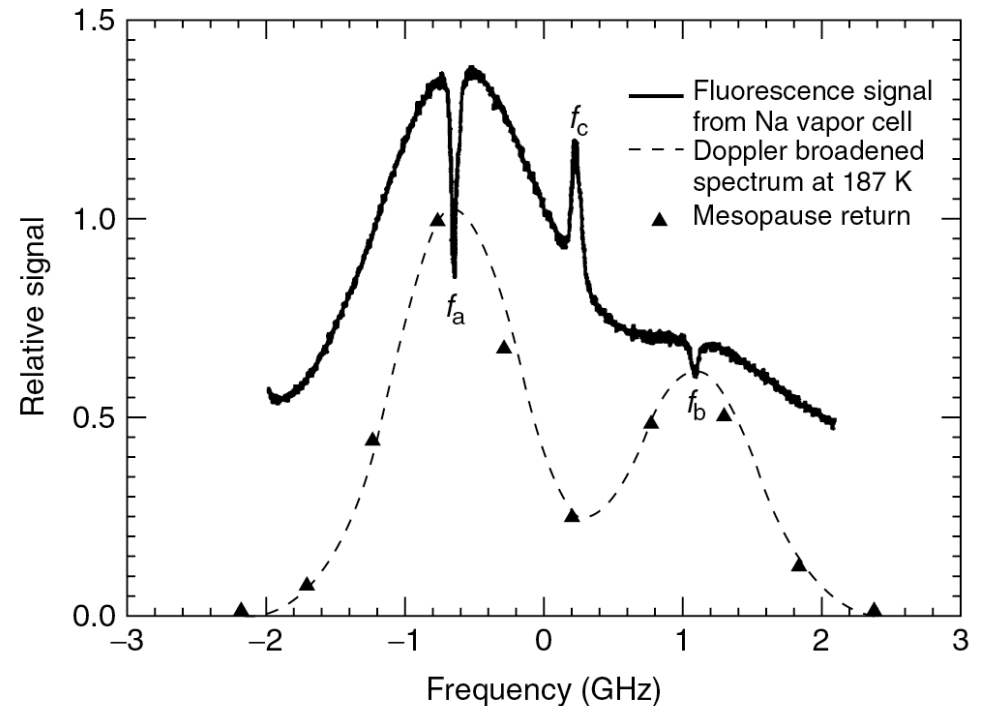
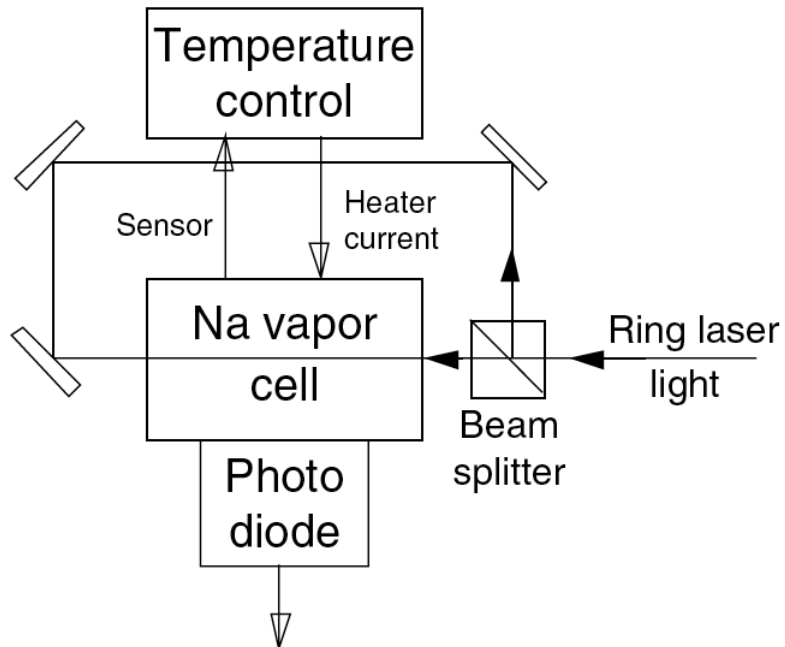
Scanning Na Lidar Results

U. Bonn LIDAR (69°N 16°E) 3. April 1984

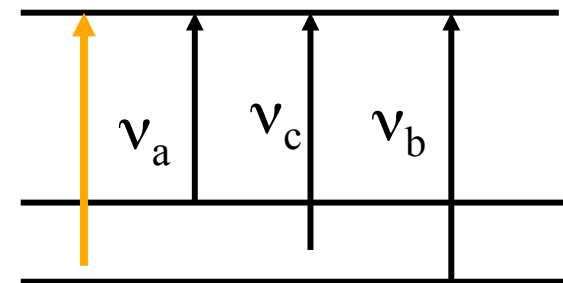




Doppler-Free Na Spectroscopy



See detailed explanation on Na Doppler-free saturation-fluorescence spectroscopy in Textbook Chapter 5.2.2.3.2

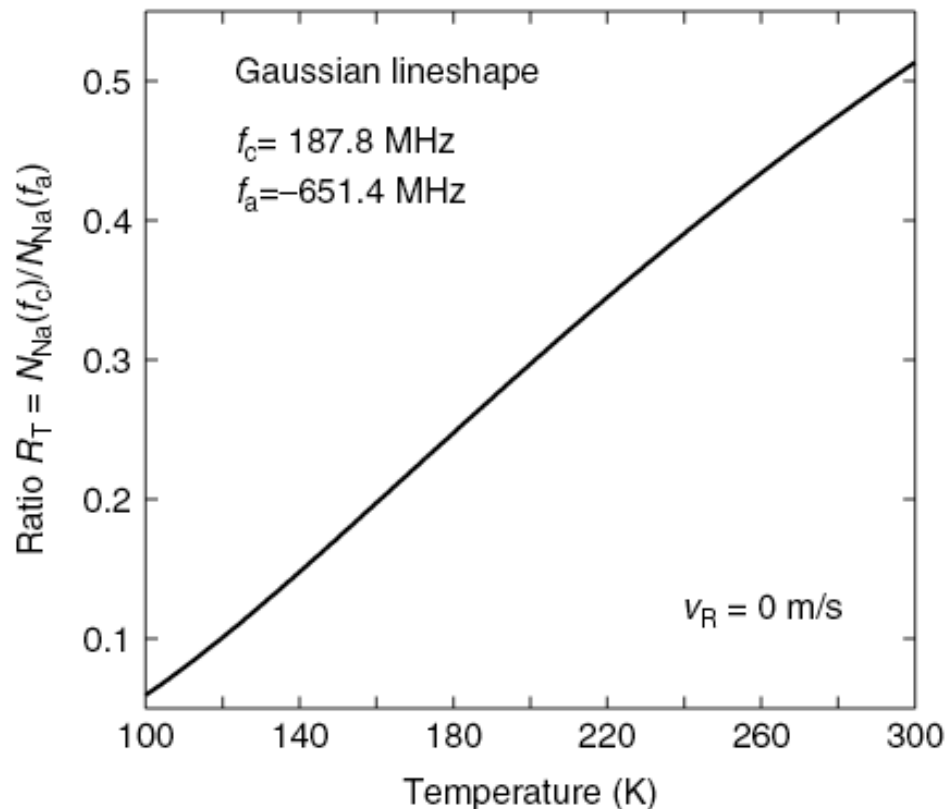


$$\nu_c = (\nu_a + \nu_b) / 2$$



Doppler Ratio Technique: 2-Frequency

$$R_T(z) = \frac{N_{norm}(f_c, z, t_1)}{N_{norm}(f_a, z, t_2)} = \frac{\sigma_{eff}(f_c, z) n_{Na}(z, t_1)}{\sigma_{eff}(f_a, z) n_{Na}(z, t_2)} \approx \frac{\sigma_{eff}(f_c, z)}{\sigma_{eff}(f_a, z)}$$



$$N_{norm}(f, z, t) = \frac{N_{Na}(f, z, t)}{N_R(f, z, t) T_c^2(f, z)}$$

$$N_{norm}(f, z, t) = \frac{\sigma_{eff}(f) n_{Na}(z)}{\sigma_R(\pi, f) n_R(z_R)} \frac{z_R^2}{4\pi z^2}$$

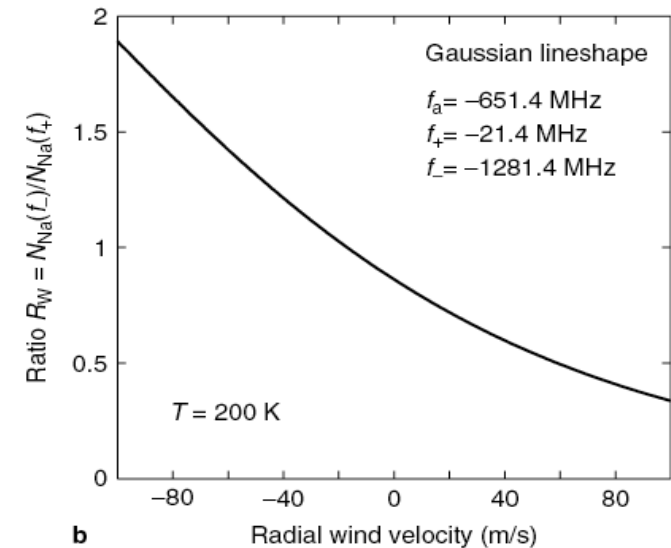
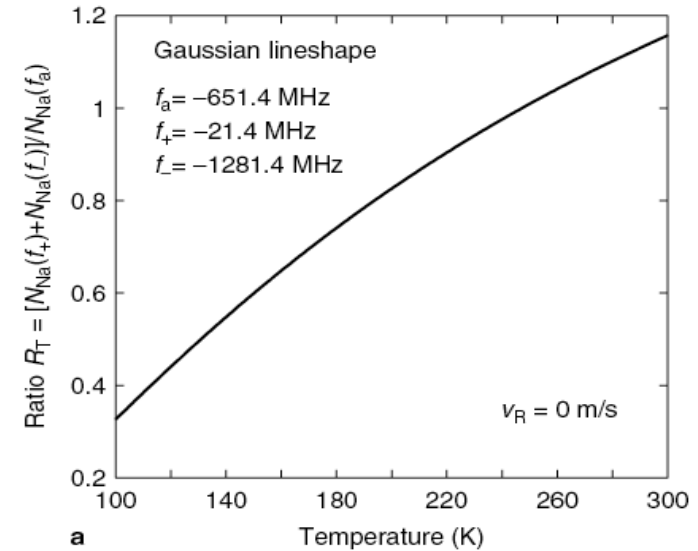


Doppler Ratio Technique: 3-Frequency

$$R_T(z) = \frac{N_{norm}(f_+, z, t_1) + N_{norm}(f_-, z, t_2)}{N_{norm}(f_a, z, t_3)}$$

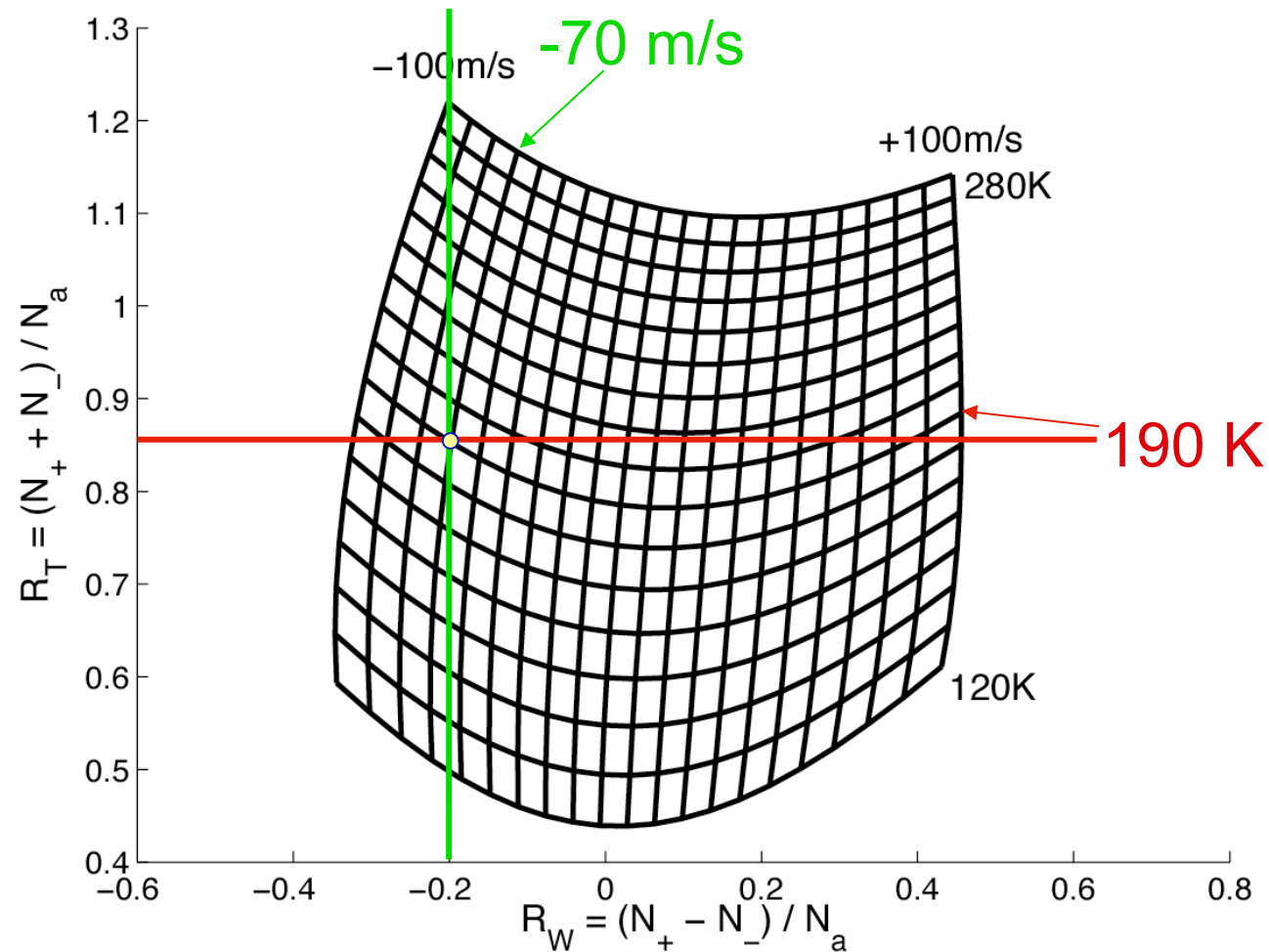
$$\approx \frac{\sigma_{eff}(f_+, z) + \sigma_{eff}(f_-, z)}{\sigma_{eff}(f_a, z)}$$

$$R_W(z) = \frac{N_{norm}(f_-, z, t_2)}{N_{norm}(f_+, z, t_1)} \approx \frac{\sigma_{eff}(f_-, z)}{\sigma_{eff}(f_+, z)}$$





Na Doppler Lidar Calibration Curves





Main Ideas Behind Ratio Technique

- ❑ Three unknown parameters (temperature, radial wind, and Na number density) require 3 lidar equations at 3 frequencies as minimum \Rightarrow highest resolution.
- ❑ In the ratio technique, Na number density is cancelled out. So we have two ratios R_T and R_W that are independent of Na density but both dependent on T and W.
- ❑ The idea is to derive temperature and radial wind from these two ratios first, and then derive Na number density using computed temperature and wind at each altitude bin.
- ❑ However, because the Na extinction coefficient is involved, the upper bins are related to lower bins, and extinction coefficient is related to Na density and effective cross-section. The solution is to start from the bottom of the Na layer and then work bin by bin to the layer top.



Summary

- ❑ The key point to measure temperature is to find and use temperature-dependent effects and phenomena to make measurements.
- ❑ Doppler technique utilizes the Doppler effect (frequency shift and linewidth broadening) by moving particles to infer wind and temperature information. It is widely applied in lidar, radar and sodar technique as well as passive optical remote sensing.
- ❑ Resonance fluorescence Doppler lidar technique applies scanning or ratio technique to infer the temperature and wind from the Doppler spectroscopy, while the Doppler spectroscopy is inferred from intensity ratio at different frequencies.