

# Optical Remote Sensing with Coherent Doppler Lidar

## Part 1: Background and Doppler Lidar Hardware

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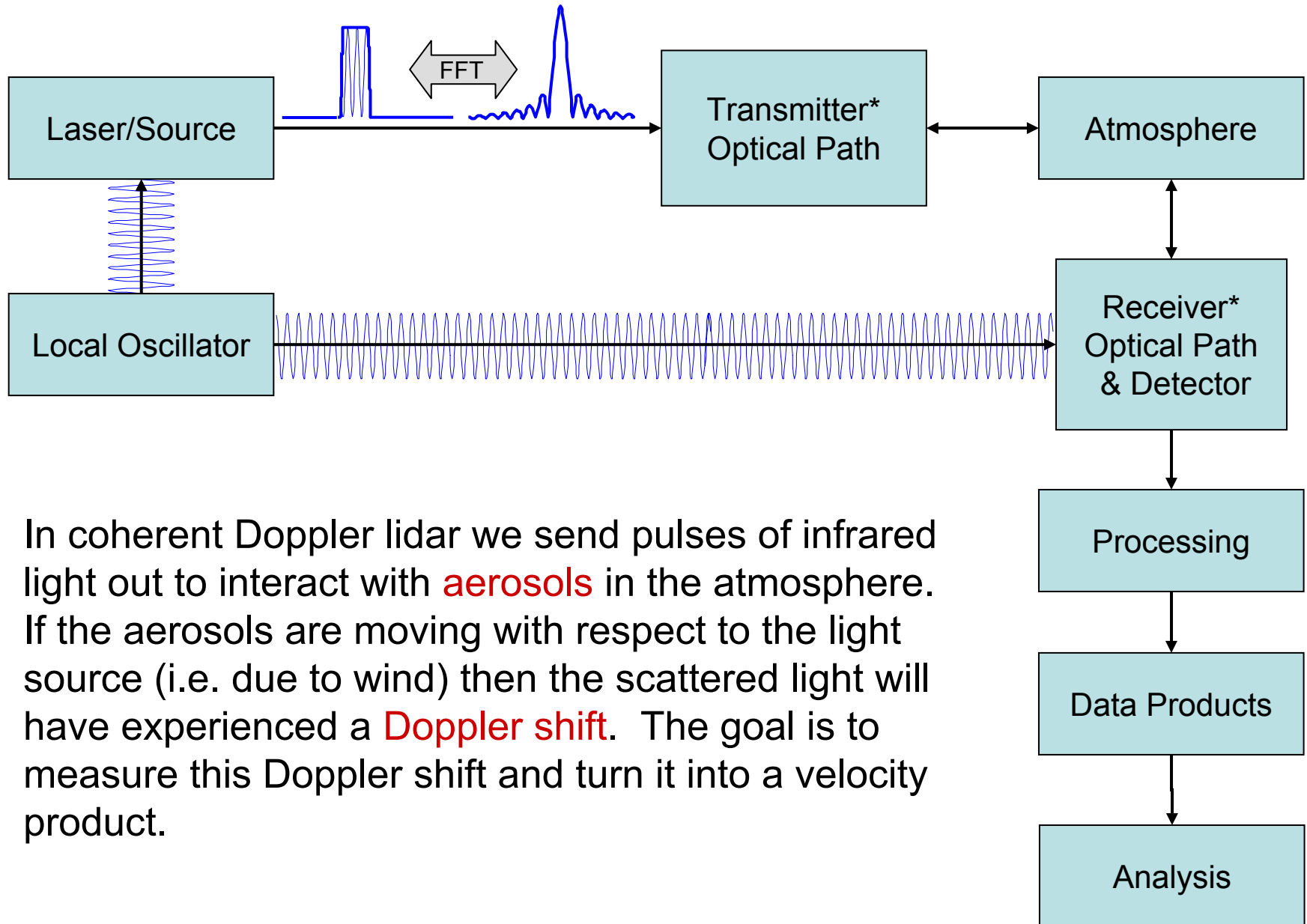
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Earth System Research Laboratory  
Chemical Sciences Division

<http://www.etl.noaa.gov>

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# Coherent Doppler Lidar



In coherent Doppler lidar we send pulses of infrared light out to interact with **aerosols** in the atmosphere. If the aerosols are moving with respect to the light source (i.e. due to wind) then the scattered light will have experienced a **Doppler shift**. The goal is to measure this Doppler shift and turn it into a velocity product.

\* Transmitter & receiver paths share some common optics

# Coherent Doppler Lidar: Return Power

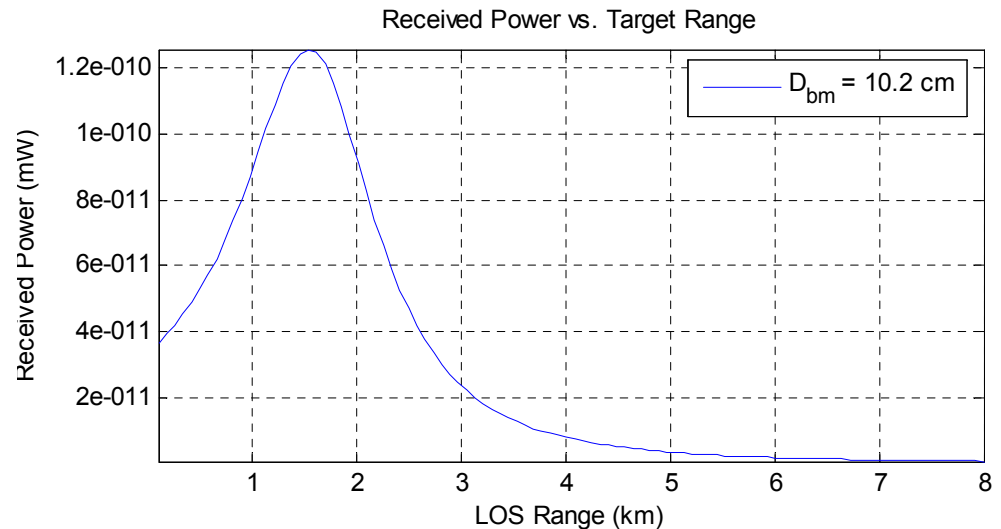
The received power,  $P_r$  is theoretically given by

$$P_r = \int_0^{\infty} \frac{A_{eff} \beta T^2}{R^2} P_T \left( \lambda, t - \frac{2R}{c} \right) dR$$

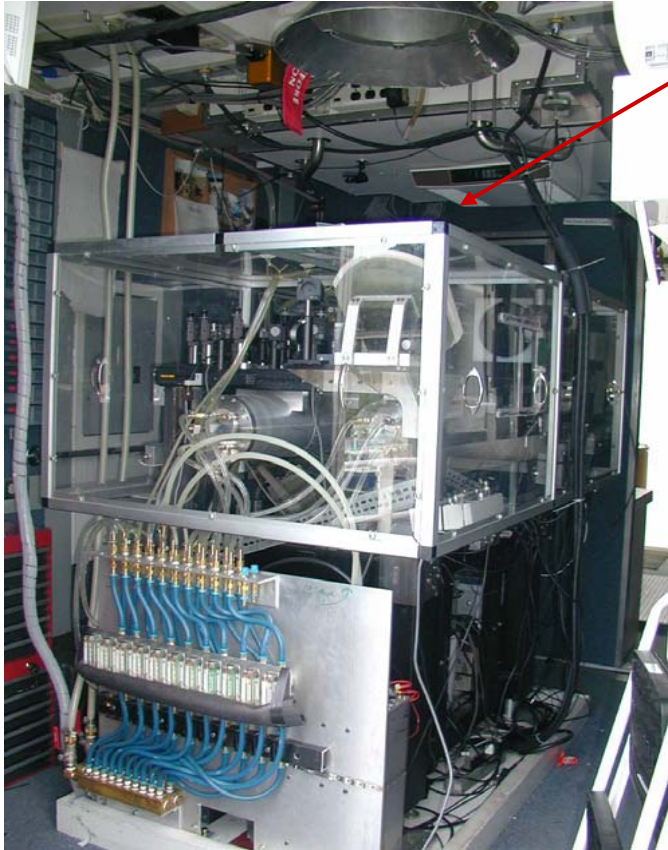
- $P_T$  = Transmitted laser power (Watts) for wavelength  $\lambda$ , range  $R$  and time  $t$ .
- $R$  = range (meters)
- $\beta$  = aerosol backscatter coefficient ( $\text{m}^{-1} \text{sr}^{-1}$ ),
- $T$  = one-way atmospheric transmission.
- $A_{eff}$  is the effective antenna area of the transceiver for a target at range  $R$ .

For aerosol targets distributed in range (relative to the pulse length) the received power at the lidar  $P_r$  can be approximated as

$$P_r = \frac{A_{eff} \beta T^2}{2R^2} c E_T$$



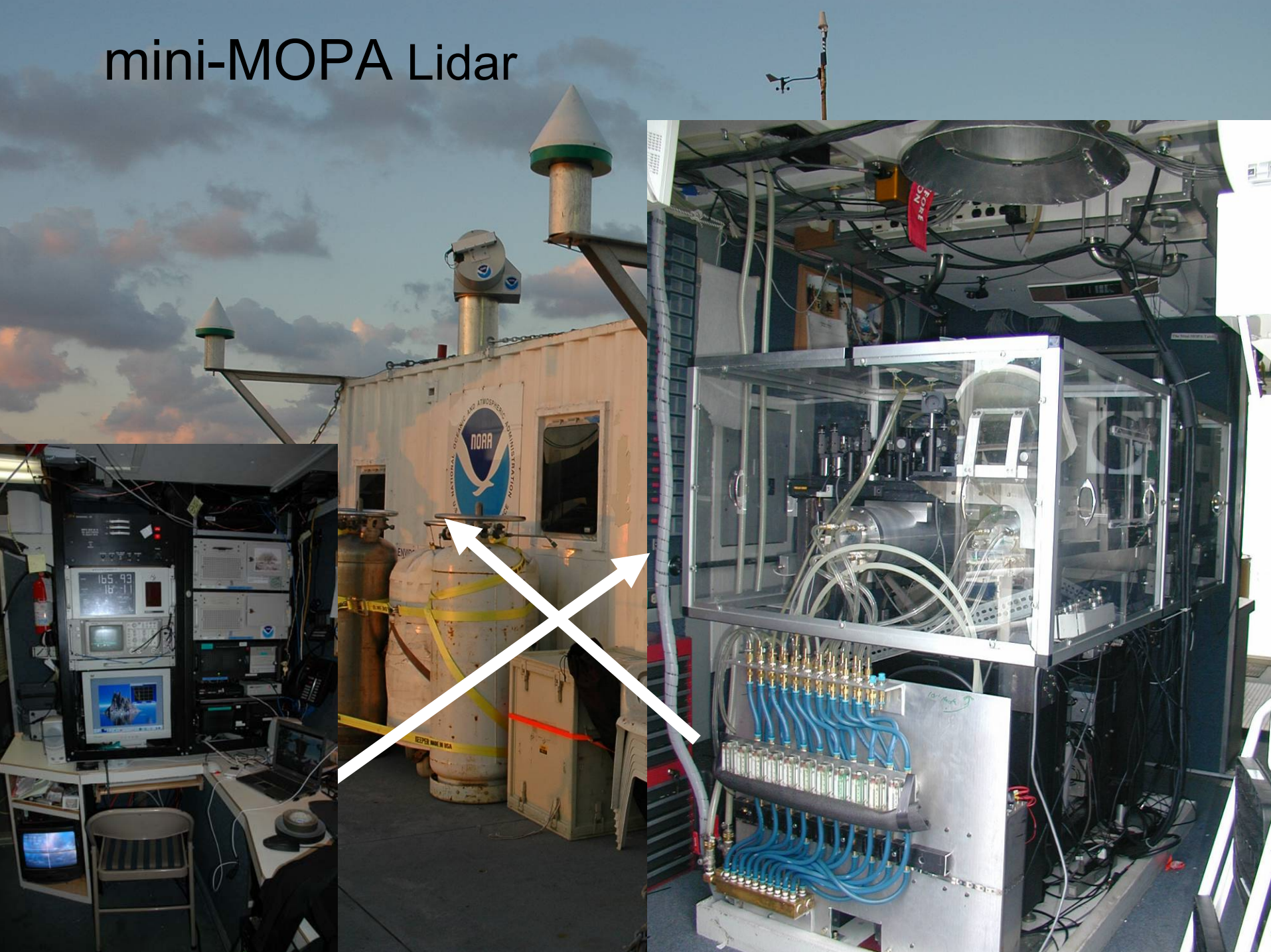
# NOAA ESRL Lidars



- Mini-MOPA
- HRDL
- OPAL
- TOPAZ
- ABDIAL
- DABUL
- Fish Lidars
- CODI
- TEAC0
- ABAEL



# mini-MOPA Lidar

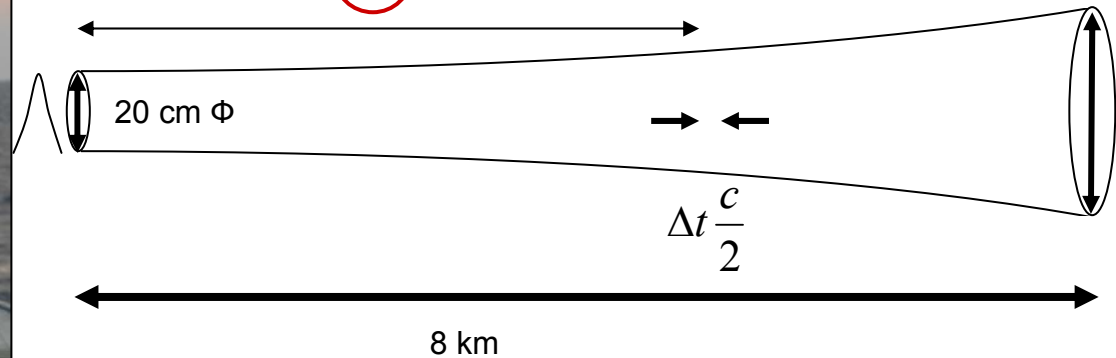


# Coherent Doppler Lidar

## Lidar measurement volume:

- Diffraction limited divergence (60  $\mu$ rad)
- “Spotlight” beam can measure to within a few meters of the surface (no side lobes)
- 30-150 m measurement volume (range resolution) along the beam (Instrument dependent)

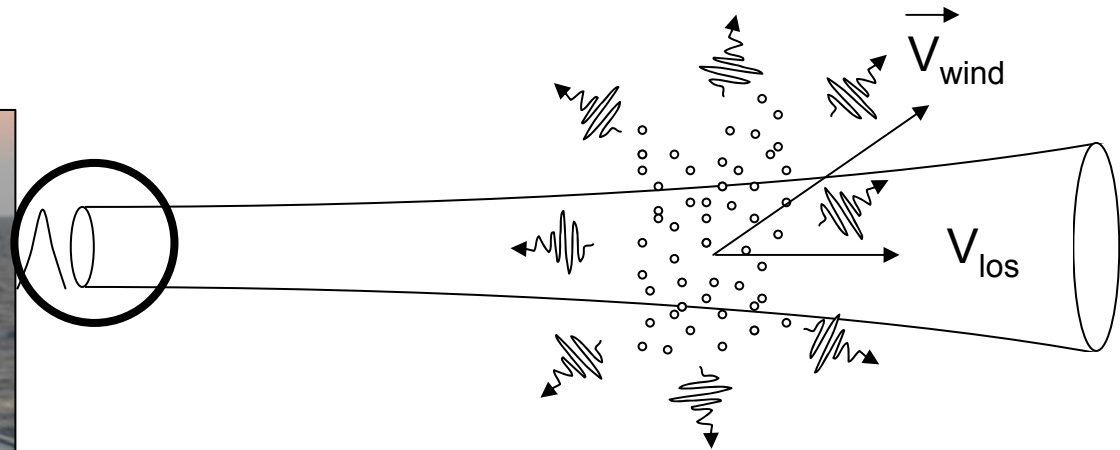
$$P_r = \frac{A_{eff} \beta T^2}{2R^2} c E_T$$



# Coherent Doppler Lidar

## Light Scattering : $\sim 2 \mu\text{m}$ & $10 \mu\text{m}$

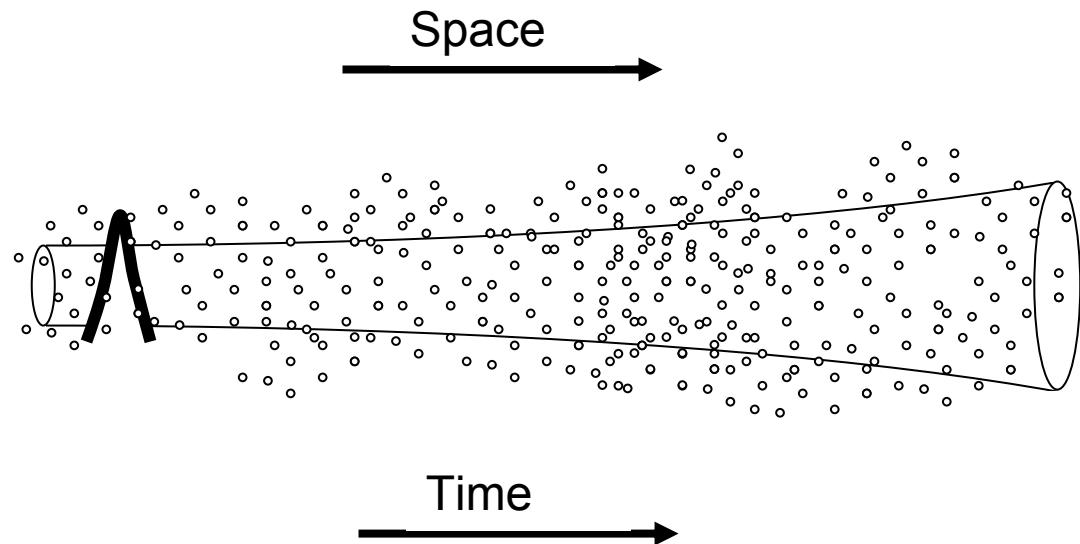
- The targets are aerosol particles
- The light scatters off the aerosol in all directions
- Part of the scattered light is detected – backscatter,  $\beta$
- The wind carries the aerosol scattering targets
- Doppler measurement is made to determine wind speed along the line of site



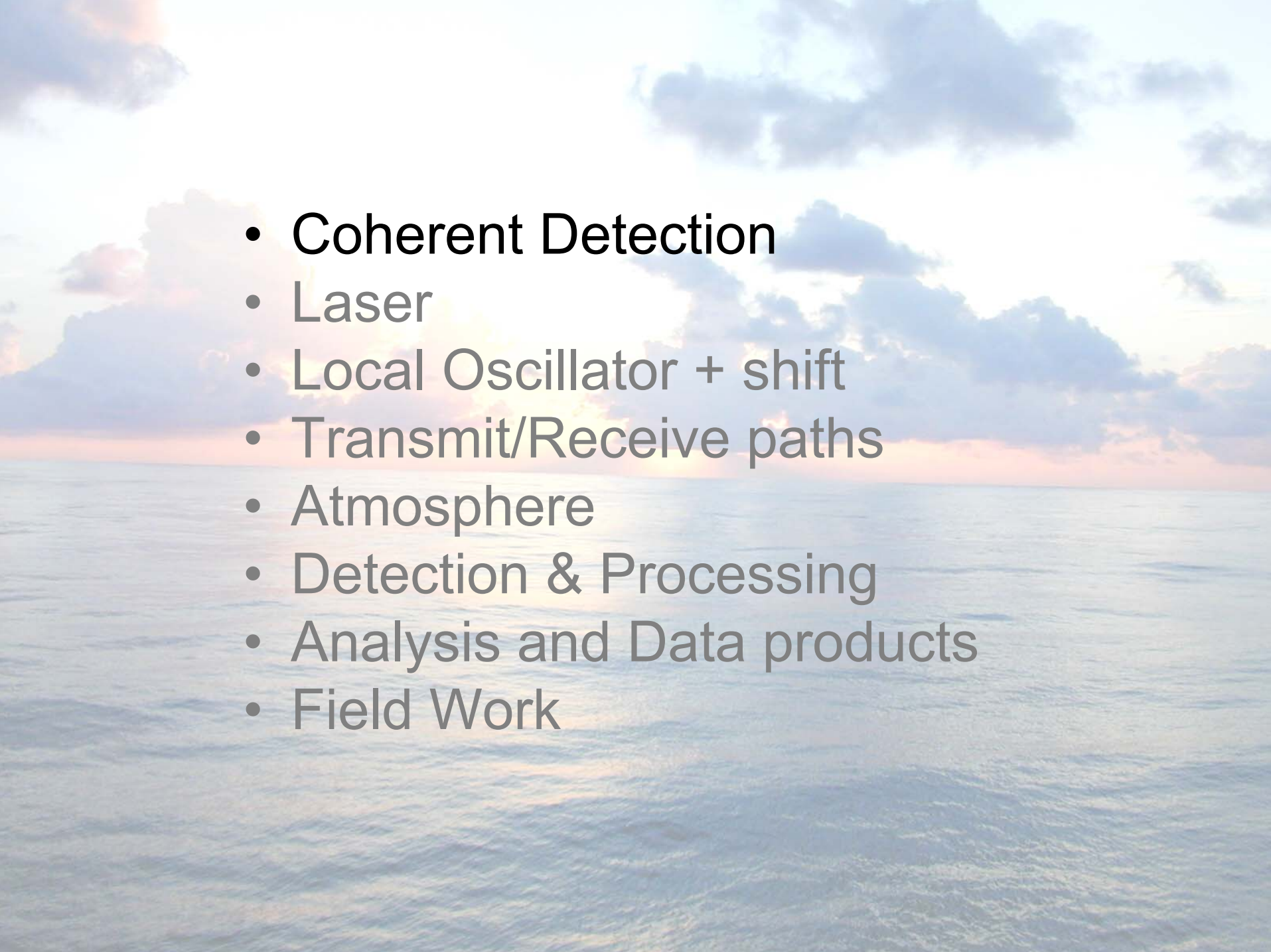
# Coherent Doppler Lidar

Light scatters from distributed target:

- For distributed aerosol
- As the pulse propagates out, a continuous signal is scattered back to the telescope and detected





- 
- **Coherent Detection**
  - Laser
  - Local Oscillator + shift
  - Transmit/Receive paths
  - Atmosphere
  - Detection & Processing
  - Analysis and Data products
  - Field Work

# Coherent Detection: The Doppler shift

- The **Doppler shift** for illumination of wavelength  $\lambda$  is given by:

$$\Delta f = \frac{2v \cos \theta_v}{\lambda} = \frac{2\nu v \cos \theta_v}{c}$$

Where  $v$  is the velocity of the aerosol(s) (e.g. wind speed) and  $\theta_v$  is the angle between the wind direction and the lidar line of sight (LOS)

For a 15 m/s wind speed, the Doppler shift for 2 $\mu$ m light ( $f_{Dopp} = 1.5 \times 10^{14}$  Hz) is 15 MHz.

- The returning illumination has a frequency of

$$f_{return} = f + f_{Dopp} = 1.50000015 \times 10^{14} \text{ Hz.}$$

- Cutoff frequencies of our detectors are around GHz.
- How can we detect such small Doppler shifts in frequencies way above detection limit?

# Coherent Detection Detecting Doppler Shifts

We can't detect the frequency of light - but we can detect the “beat” (i.e. difference) signal between two light beams of slightly different frequency...

So, we create two beams: a **local oscillator** (LO) and a **power oscillator** (PO). The Local oscillator has frequency  $f_{LO}$ .

We make sure that the PO has a known frequency offset (i.e.  $f_{offset} = 10$  MHz, 100 MHz) from that of the LO, or  $f_{PO} = f_{LO} + f_{offset}$ .

This PO beam goes out into the atmosphere. The light that returns (scattering off of aerosols) may have been Doppler shifted by  $f_{Dopp}$  for a total frequency offset of

$$f_a = f_{Dopp} + f_{offset} + f_{LO}$$

# Coherent Detection

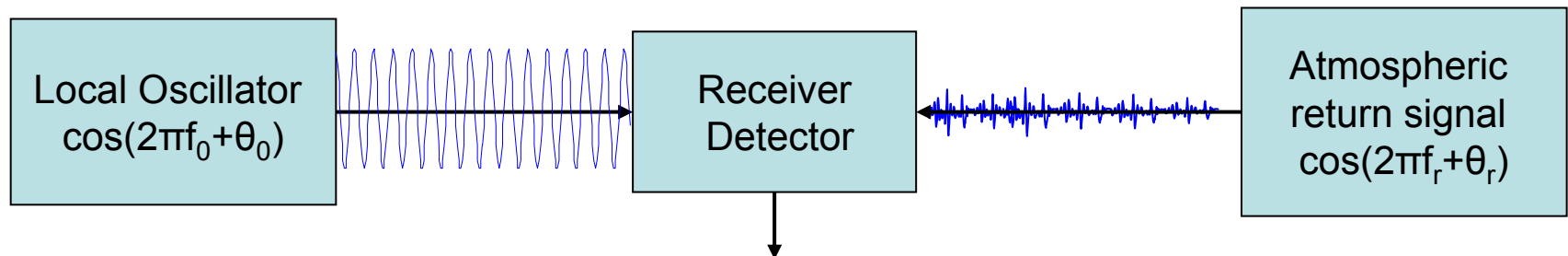
The atmospheric return signal and the signal from the local oscillator are both incident on the detector.

Their electric fields add to create the total electric field incident on the detector:

$$E_a = A_a \cos(j2\pi f_a t + \varphi_a)$$

$$E_{LO} = A_{LO} \cos(j2\pi f_{LO} t + \varphi_{LO})$$

$$E_{tot} = A_a \cos(j2\pi f_a t + \varphi_a) + A_{LO} \cos(j2\pi f_{LO} t + \varphi_{LO})$$



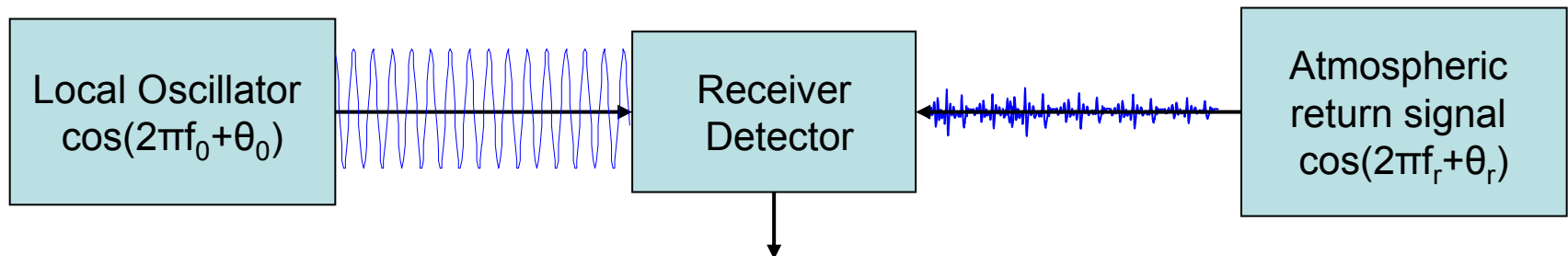
# Coherent Detection

The detector actually “sees” optical power or:

$$\begin{aligned} |E_{tot}|^2 &= |A_a \cos(j2\pi f_a t + \varphi_a) + A_{LO} \cos(j2\pi f_{LO} t + \varphi_{LO})|^2 \\ &= A_a^2 |\cos(j2\pi f_a t + \varphi_a)|^2 + A_{LO}^2 |\cos(j2\pi f_{LO} t + \varphi_{LO})|^2 \\ &\quad + 2A_a A_{LO} \cos(j2\pi f_a t + \varphi_a) \cos(j2\pi f_{LO} t + \varphi_{LO}) \end{aligned}$$

The product of cosines leads to a sum and a difference:

$$\begin{aligned} |E_{tot}|^2 &= A_a^2 |\cos(j2\pi f_a t + \varphi_a)|^2 + A_{LO}^2 |\cos(j2\pi f_{LO} t + \varphi_{LO})|^2 \\ &\quad + 2A_a A_{LO} \cos(j2\pi(f_a + f_{LO})t + (\varphi_a + \varphi_{LO})) \\ &\quad + 2A_a A_{LO} \cos(j2\pi(f_a - f_{LO})t + (\varphi_a - \varphi_{LO})) \end{aligned}$$

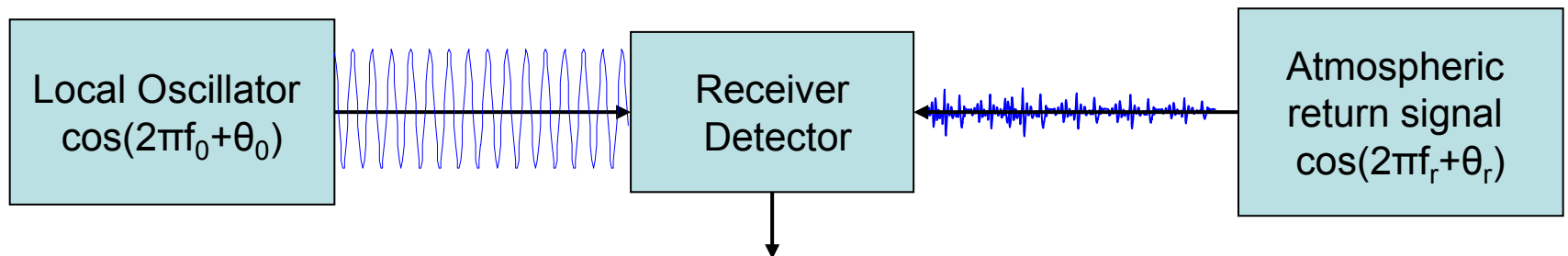


# Coherent Detection

The high frequency (i.e. the sum of LO and atmospheric frequencies) is too high to detect. The other terms contribute to a DC offset, and the difference frequency is what gives us our signal:

$$|E_{tot}|^2 = |E_a|^2 + |E_{LO}|^2 + A_a A_{LO} \cos(j2\pi(f_a - f_{LO})t + (\varphi_a - \varphi_{LO}))$$

In terms of power - the optical power on the detector is given by:

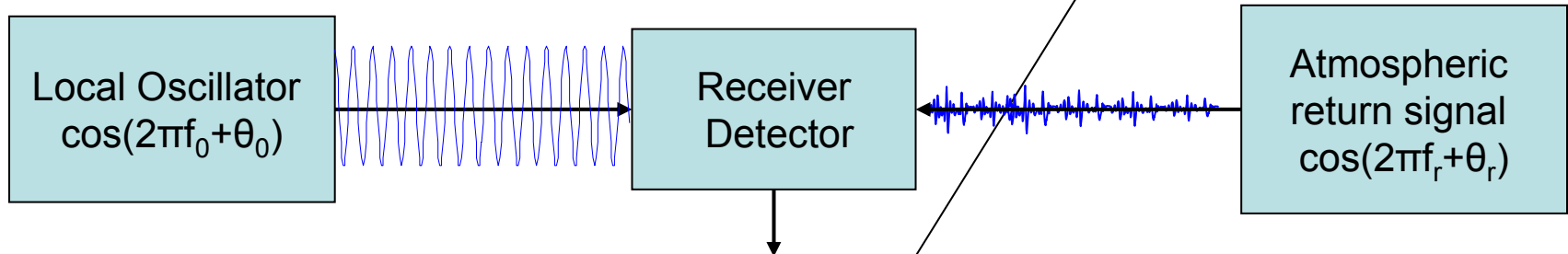


$$P_{sig} = P_a + P_{LO} + 2\sqrt{P_a P_{LO}} \cos(j2\pi(f_a - f_{LO})t + (\varphi_a - \varphi_{LO}))$$

# Coherent Detection

The detector current is then given by:

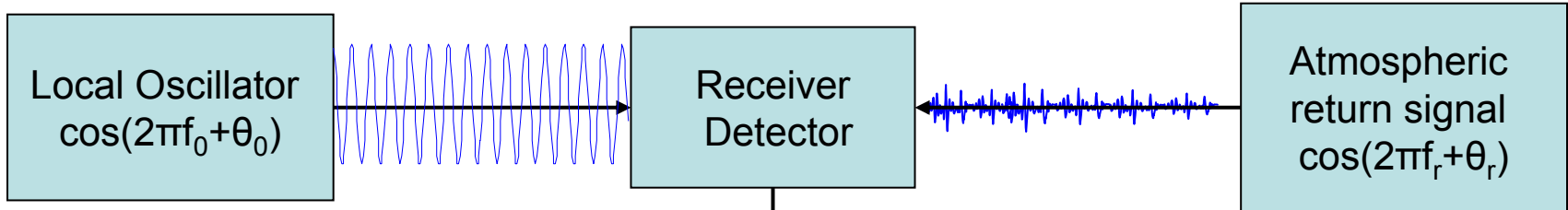
$$i_{sig} = \left( \frac{\eta e P_{sig}}{h\nu} \right) = i_a + i_{LO} + 2\sqrt{i_a i_{LO}} \cos(j2\pi(f_a - f_{LO})t + (\varphi_a - \varphi_{LO}))$$



Remember  $f_a - f_{LO} = f_{Dopp} + f_{offset} \sim \text{Mhz}$

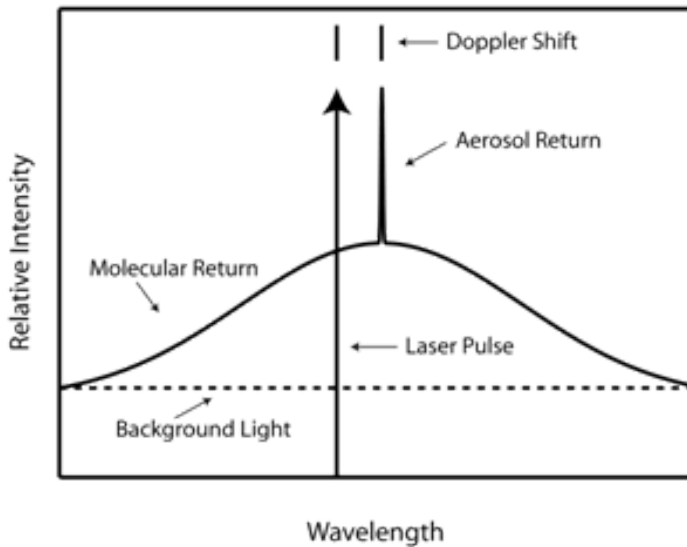
We know  $f_{offset}$ ...so we can find the Doppler shift frequency.

# Coherent Detection



$$f_{\text{detected}} = f_a - f_{LO} = f_{Dopp} + f_{\text{offset}} \sim \text{Mhz}$$

We assume that  $f_{LO}$  is the same at 20+km (or 66.7  $\mu\text{s}$  – at least) as it was when we sent the pulse out – Not always true for UV sources



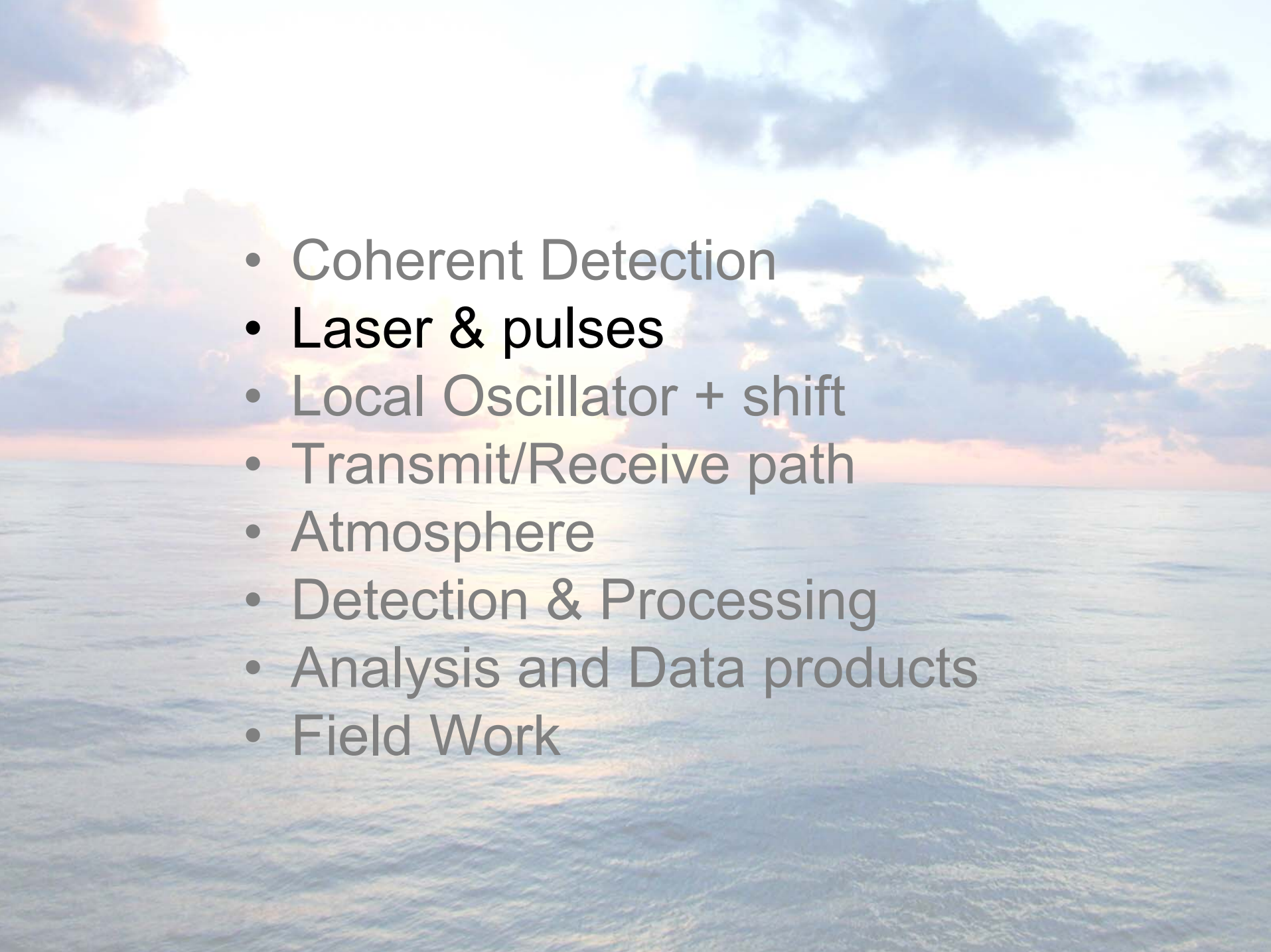
Also consider the spread of frequencies in the return signal –  $f_{Dopp}$  is **not** a single frequency.

Rayleigh vs. Mie scattering



# IR vs. UV in heterodyne detection

Property	IR	UV
Linewidth/ Temporal Coherence	kHz → 10s of km and longer (100 km)	Old: GHz → meters New: MHz → 100s m
Scattering/BW	Mie – pulse transform limited	Rayleigh (very wide) & Mie
Detection noise	Shot noise limited by LO	LO Shot noise + Rayleigh scattering
Aerosol sampling BW (SNR $\propto 1/\text{BW}$ ) $\Delta f = \frac{2v}{\lambda}$	2 $\mu\text{m}$ : 25 m/s needs 50 Mhz BW	355nm: 25 m/s needs ~300 Mhz
Refractive Turbulence	Some effect (less for longer $\lambda$ )	Stronger effect (less spatial coherence)

- 
- Coherent Detection
  - **Laser & pulses**
  - Local Oscillator + shift
  - Transmit/Receive path
  - Atmosphere
  - Detection & Processing
  - Analysis and Data products
  - Field Work

# Laser & Pulses

## Laser/Transmitter Requirements

- Narrow bandwidth (i.e.  $\sim 1$  Mhz)
- Q-switched or modulated
- Low atmospheric absorption
- High pulse repetition frequency (PRF)
- 1-8 mJ per pulse
- Eyesafe

Tradeoffs between:

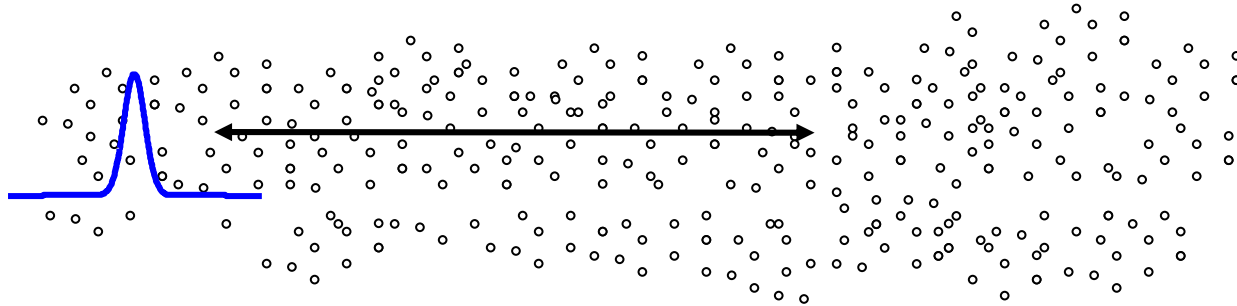
- short pulses
- pulse bandwidth
- PRF
- average power

A fun intro to lasers....

<http://www.colorado.edu/physics/2000/lasers/index.html>

# Laser & Pulses

## Time-bandwidth tradeoffs



“short” pulse

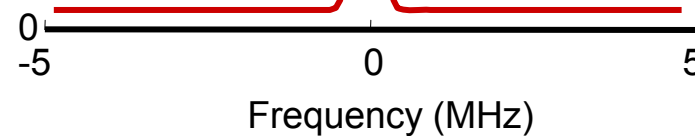
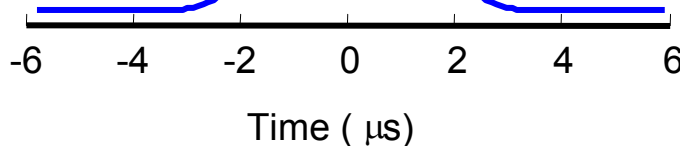
Precise in  
time/range

Ambiguous  
in frequency

“long” pulse

Ambiguous  
in time/range

Precise  
in frequency

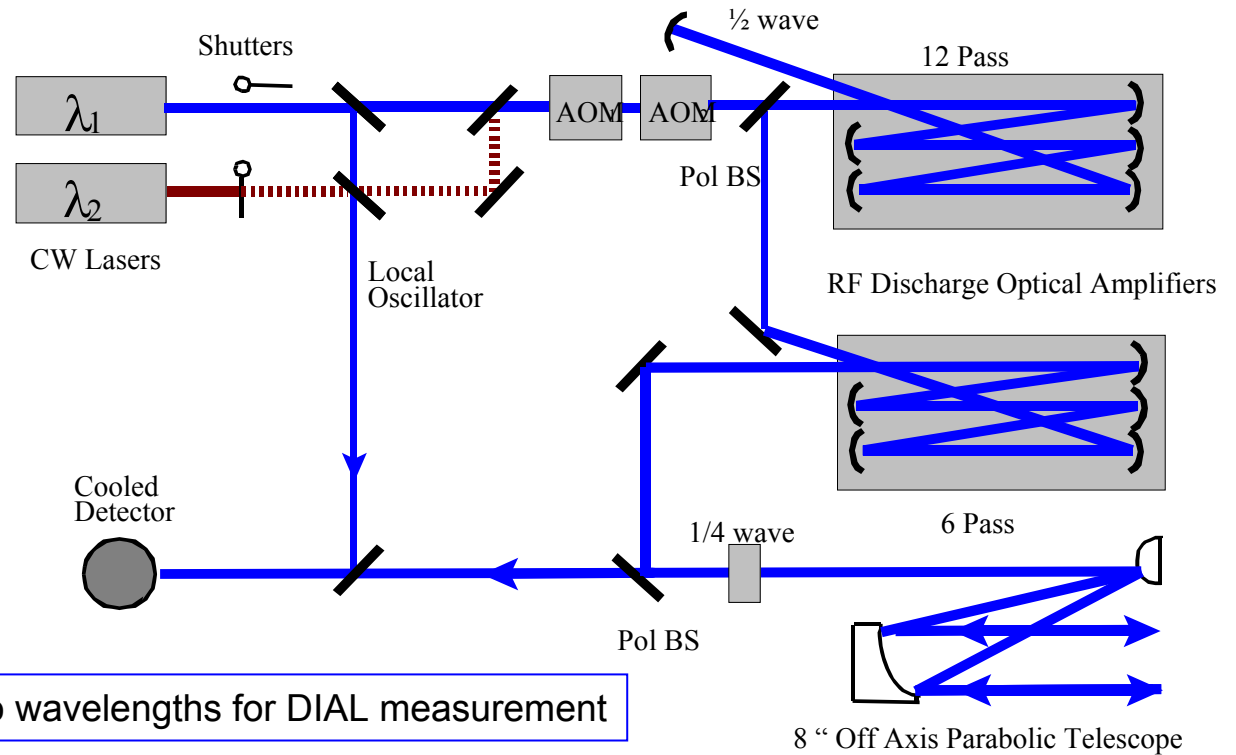


How are the pulses created?

# Laser & Pulses

## Mini-MOPA

(master-oscillator/  
power-amplifier)



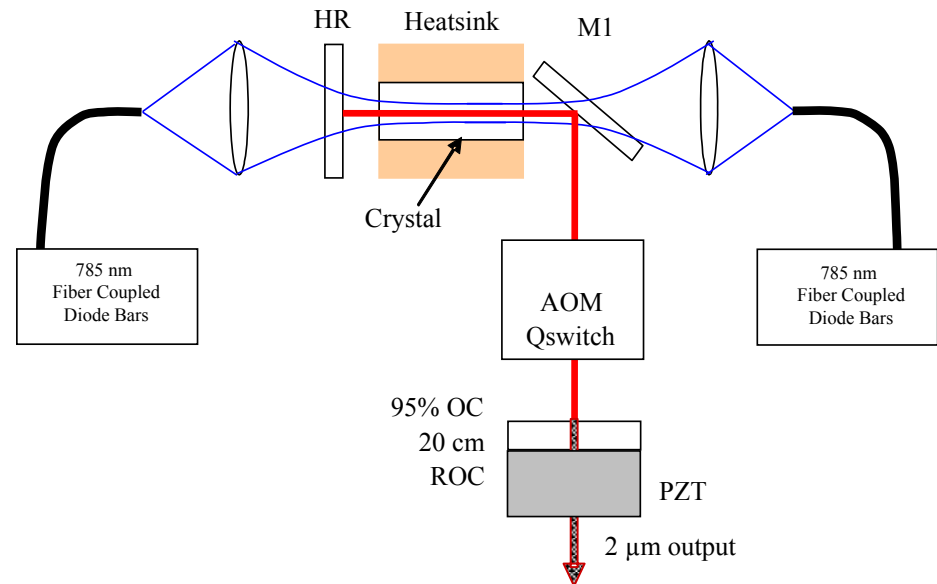
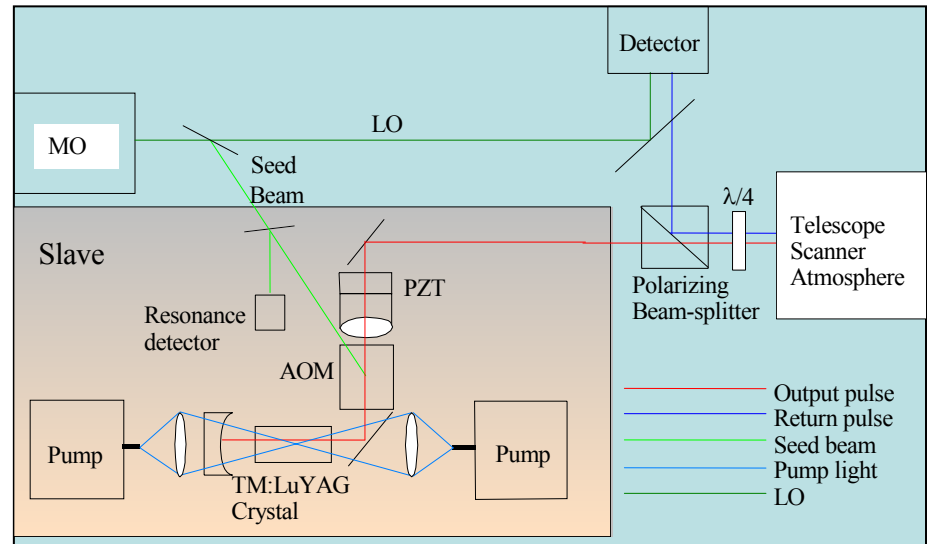
Can also alternate between two wavelengths for DIAL measurement


Wavelength	9-11 micron
Pulse Energy	0.5-2 mJ
PRF	300 Hz
Max Range	18 km
Range Resolution	45-300 m
Scanning	Full Hemispheric
Precision	10 cm/s



# Laser & Pulses: High Resolution Doppler Lidar (HRDL)

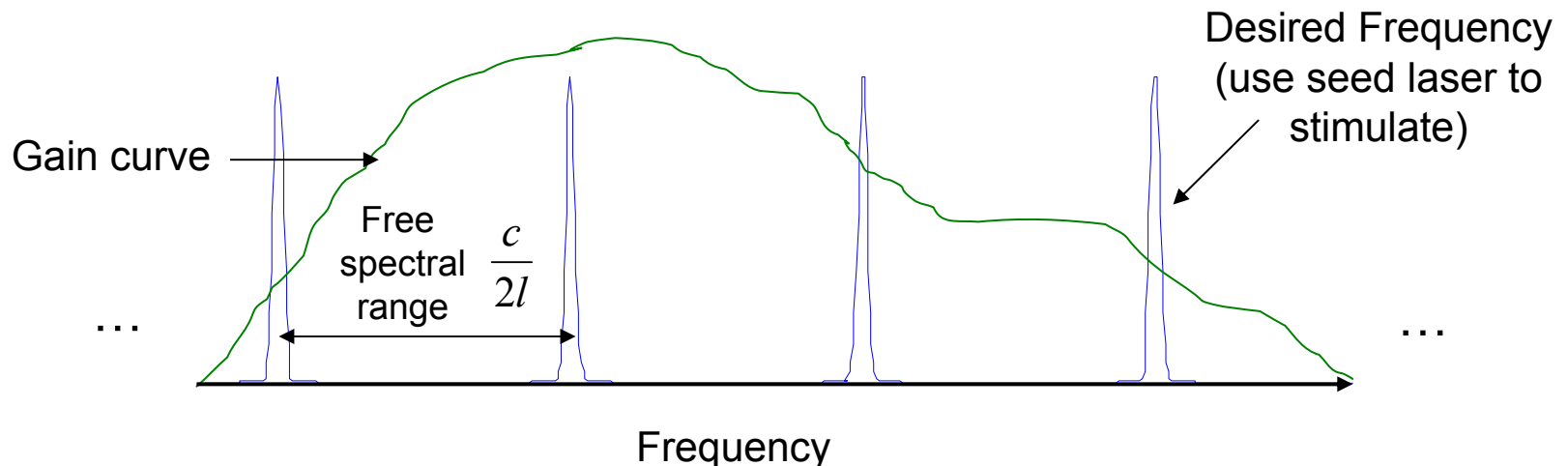
Wavelength	2.02 micron
Pulse Energy	2 mJ
PRF	200 Hz
Max Range	3-8 km
Range Res.	30 m
Beam rate	2 Hz
Scanning	Full Hemispheric
Precision	10 cm/s



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  - Laser & pulses
  - **Local Oscillator + shift**
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# Local Oscillator + shift: LO Requirements

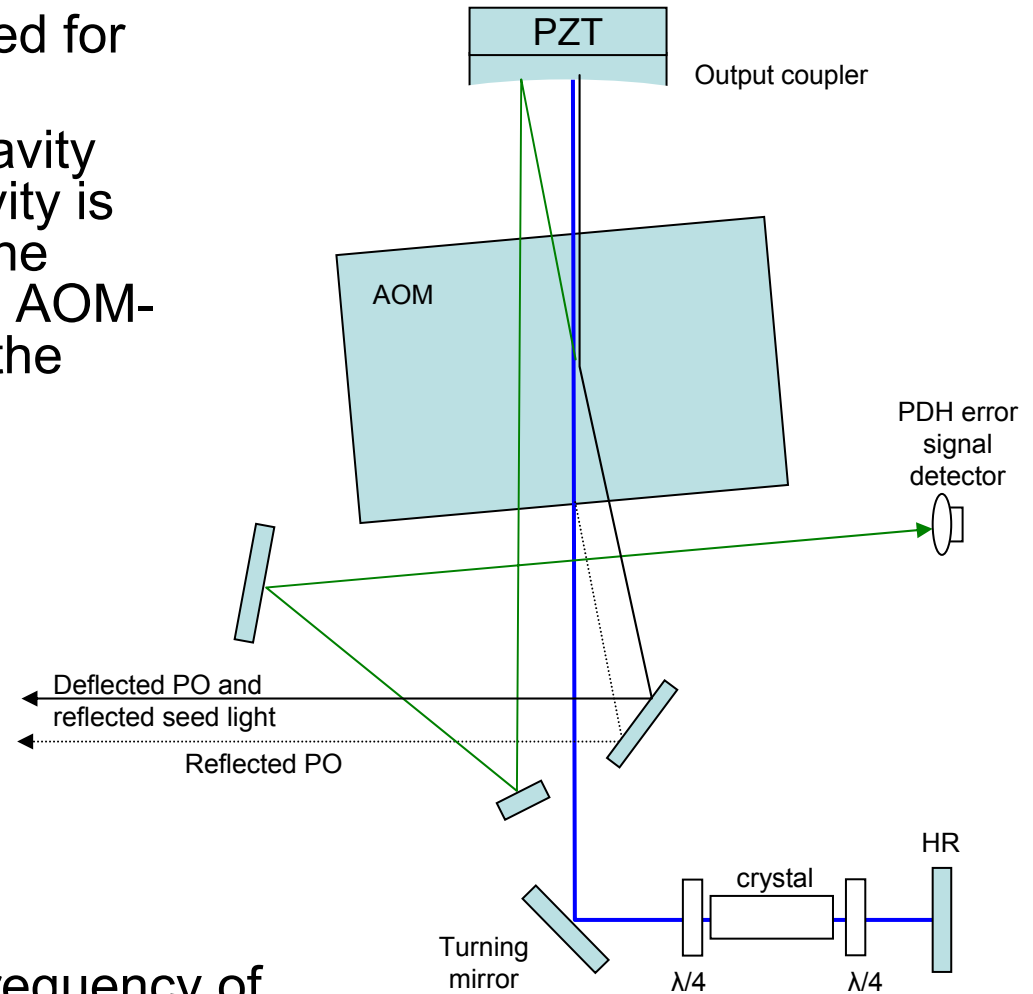
- Continuous wave – always available for heterodyne detection of return pulses from the atmosphere.
- Stable – especially over pulse separation times.
- Need a way to shift the frequency of the pulses relative to the LO (or the other way around) – we use AOMs for this.
- Sometimes the same source as the PO – sometimes a seed for the PO.

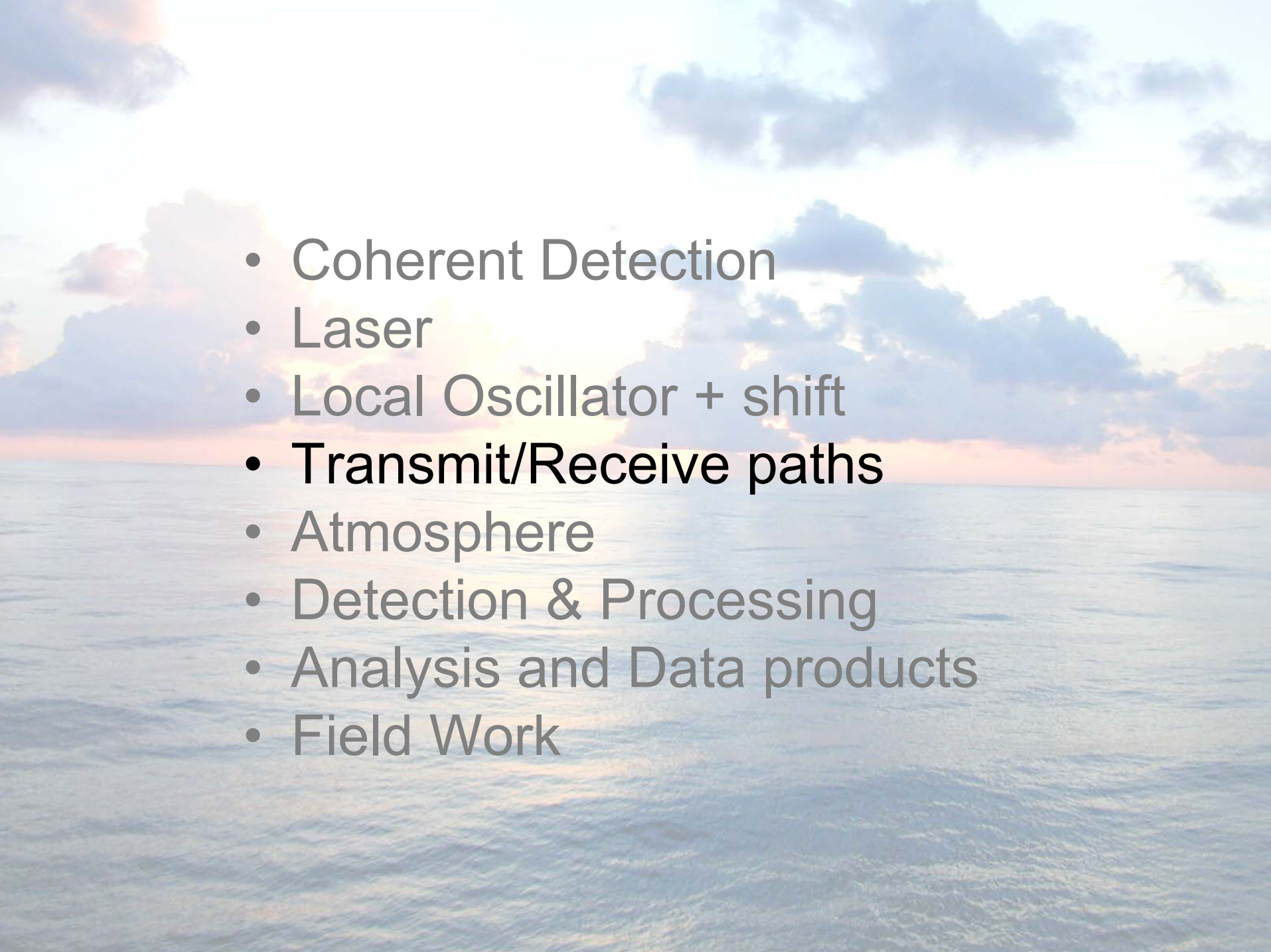




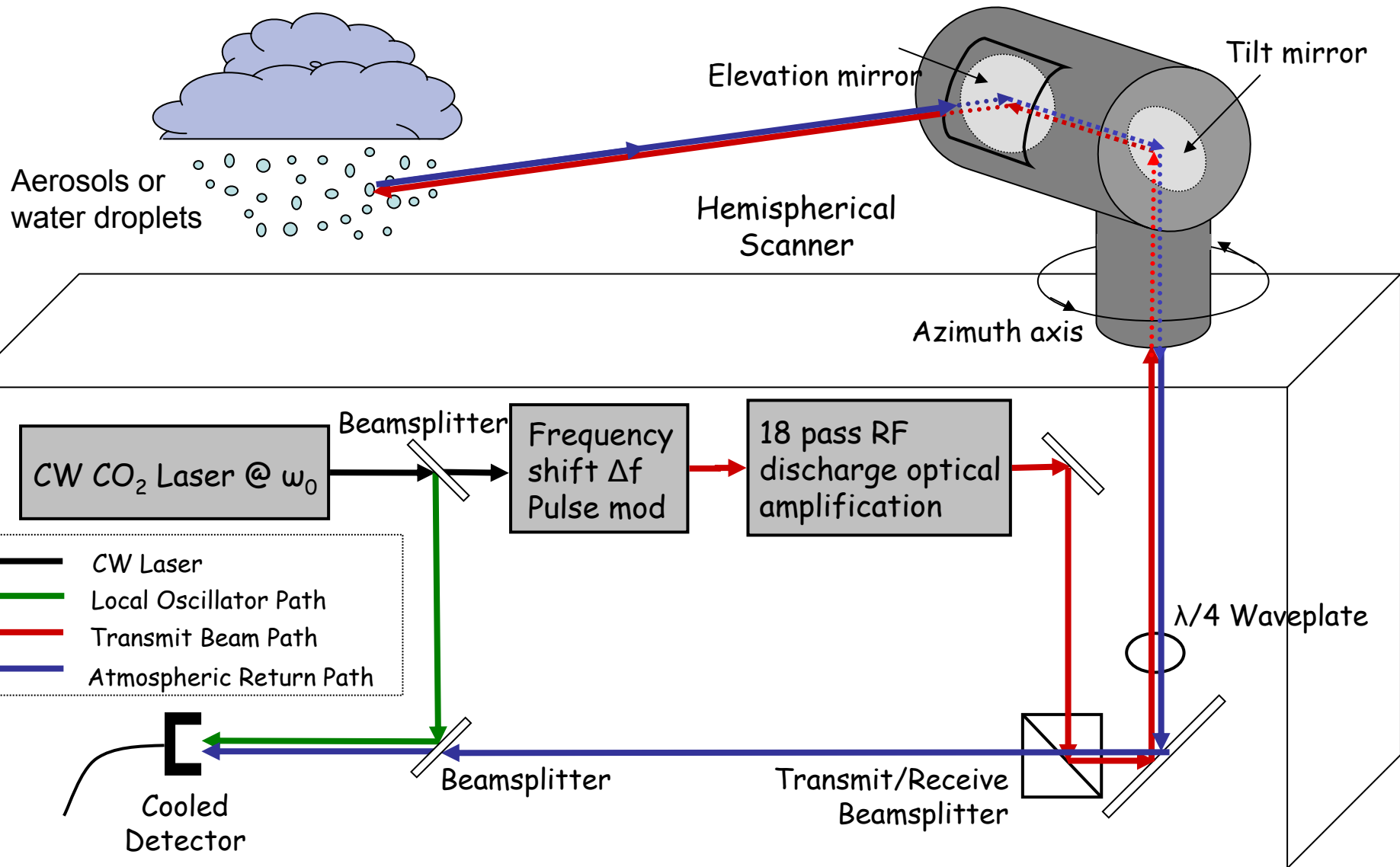
# Local Oscillator & Seed: HRDL

- The LO is a separate laser seed for the PO
- The LO is “injected” into the cavity using the AOM angle. The cavity is then adjusted to optimize for the frequency of the LO PLUS the AOM-induced frequency offset and the AOM is turned off.
- At this time, the PO light in the cavity has already started the stimulated emission process – now all the photons emit at the same frequency and phase – and the pulse is formed.
- The AOM causes the center frequency of the pulse to be 100 Mhz higher than the LO seed light.

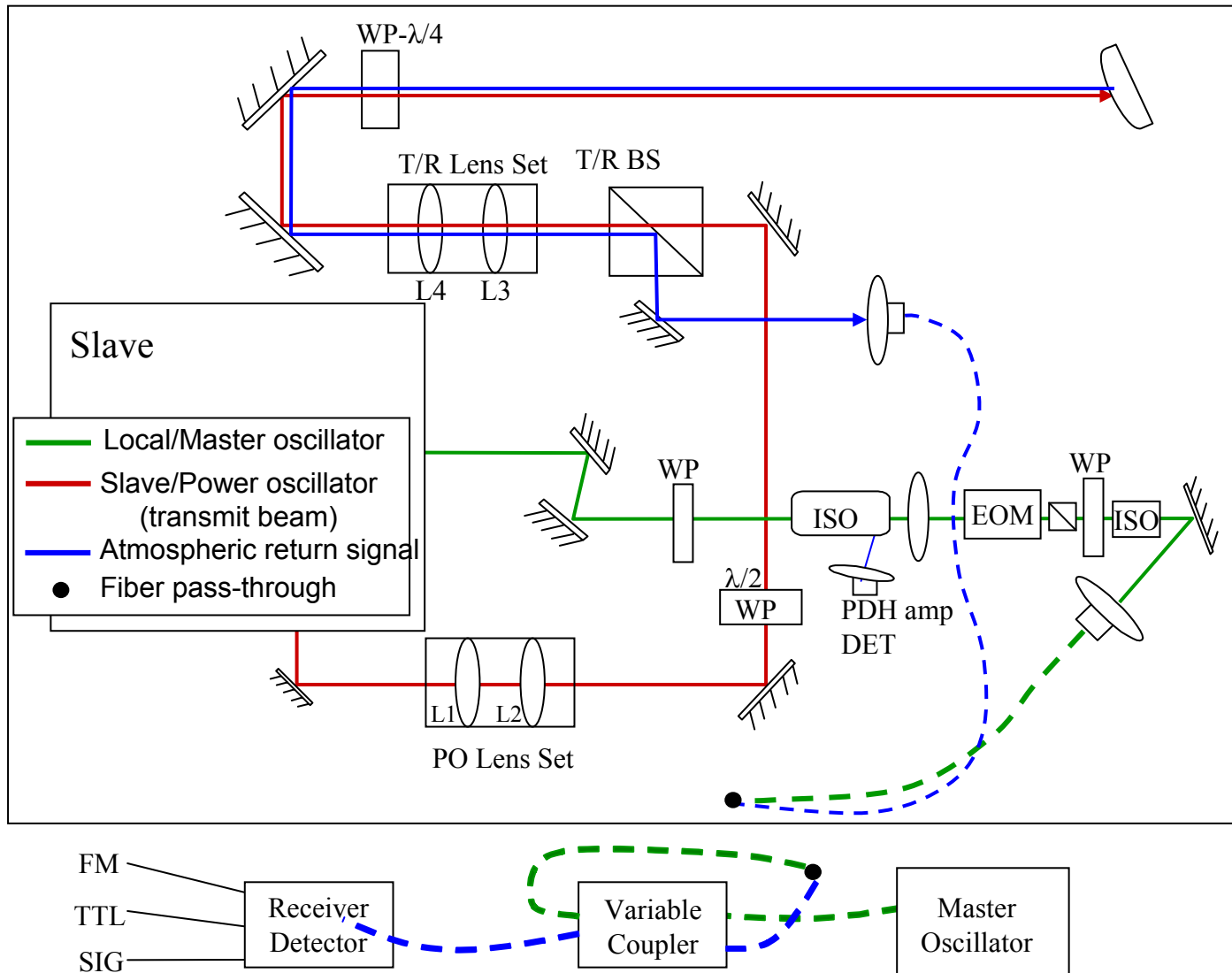



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# Mini-MOPA (master-oscillator/power-amplifier) system



# High Resolution Doppler Lidar (HRDL) system

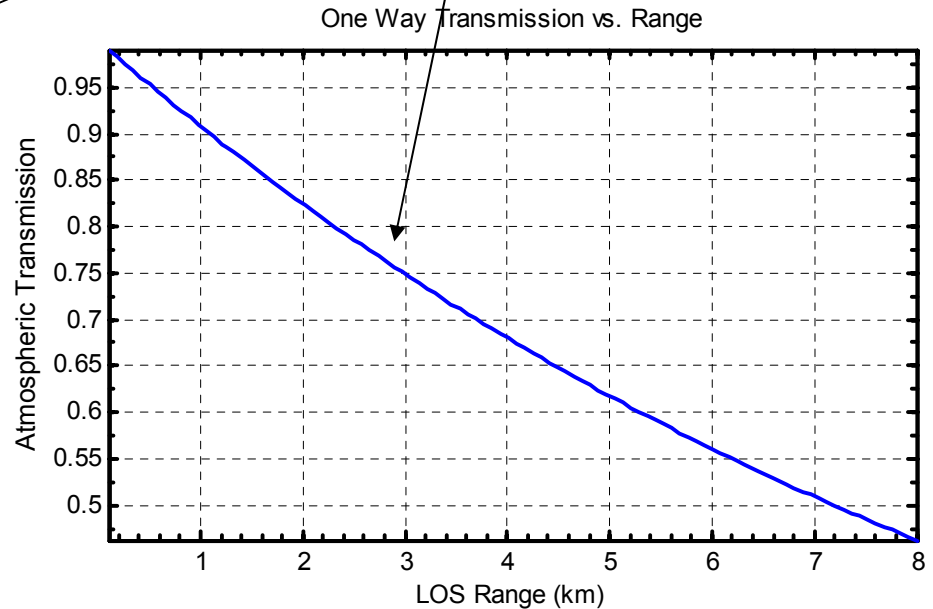
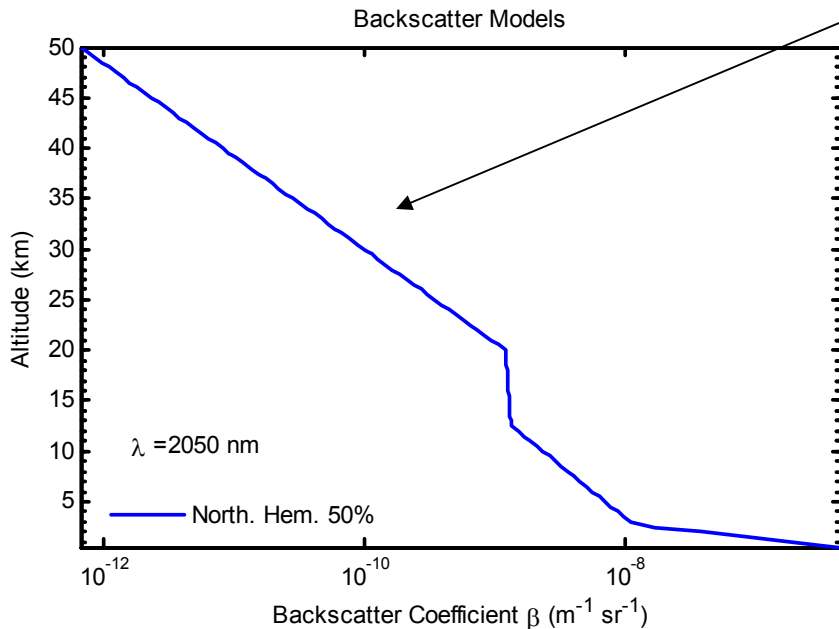


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# Atmospheric Return

- Continuous return from distributed target
- Atmosphere affects the amount of return signal according to the amount of aerosols (backscatter), extinction, and turbulence.

$$P_r = \frac{A_{eff} \beta T^2}{2R^2} cE_T$$



## The Coherent Doppler Lidar Equation

The carrier-to-noise ratio (CNR) is found using the following equation:

$$CNR = \frac{\langle |i_{het}|^2 \rangle}{\langle |i_N|^2 \rangle} = \frac{\eta P_r}{h \nu B}$$

- where  $\eta$  is an efficiency factor (less than or equal to unity) describing the noise sources in the photo-detector signal as well as optical efficiencies,
- $h$  is Plank's constant ( $6.626 \times 10^{-34}$  Joule-sec)
- $\nu$  is the optical frequency (Hz.)
- $B$  is the receiver bandwidth determined by the receiver electronics.
  - In HRDL's case,  $B$  is 50 MHz. - In MOPA's case,  $B$  is 10 MHz
- *Rule of thumb:* We need about one coherent photon per inverse BW to get 0 dB CNR – i.e. Coherent Doppler Lidar is quite sensitive.

## The Coherent Doppler Lidar Equation, cont'd

The received power,  $P_r$  is theoretically given by

$$P_r = \int_0^{\infty} \frac{A_{eff} \beta T^2}{R^2} P_T \left( \lambda, t - \frac{2R}{c} \right) dr$$

$P_T$  = Transmitted laser power (Watts) for wavelength  $\lambda$ , range  $R$  and time  $t$ ,

- $R$  = range (meters)
- $\beta$  = aerosol backscatter coefficient ( $\text{m}^{-1} \text{sr}^{-1}$ ),
- $T$  = one-way atmospheric transmission.
- $A_{eff}$  is the effective antenna area of the transceiver for a target at range  $R$ .

For aerosol targets distributed in range (relative to the pulse length) the received power at the lidar  $P_r$  can be approximated as

$$P_r = \frac{A_{eff} \beta T^2}{2R^2} cE_T$$



## The Coherent Doppler Lidar Equation, cont'd

The effective area is effected by the Gaussian beam expansion and transmitter focus parameters as well as turbulence and is given by

$$\frac{1}{\langle A_{eff} \rangle} = 2 \left( \frac{1}{A_{TR}} + \frac{1}{A_{turb}} \right)$$

Where  $A_{turb}$  is the coherence area defined by  $\pi\rho_0$ .

$A_{TR}$  is the transmit/receive area defined by

$$\frac{1}{A_{TR}} = \frac{2}{\pi D^2} + \frac{\pi D^2}{8\lambda^2} \left( \frac{1}{F} - \frac{1}{R} \right)^2$$

$D_b$  is the transmitted,  $1/e^2$  intensity, untruncated, Gaussian beam diameter in meters, F is the focus of the transmitter optics.

Thus  $A_{eff}$  is defined by

$$A_{eff} = \frac{\pi D^2}{4} \left[ 1 + \left( \frac{\pi D^2}{4\lambda R} \right)^2 \left( 1 - \frac{R}{F} \right)^2 + \frac{D^2}{2\rho_0^2} \right]^{-1}$$

## The Coherent Doppler Lidar Equation, cont'd

The turbulence parameter  $\rho_o$  is given by

$$\rho_o = \left[ 1.45k^2 \int_0^{\infty} C_n^2(R') \left(1 - \frac{R'}{R}\right)^{\frac{5}{3}} dR' \right]^{-\frac{3}{5}}$$

For constant refractive turbulence ( $C_n^2$ ) level, The above equation reduces to

$$\rho_o = \left[ 1.45k^2 C_n^2 \frac{3}{8} R \right]^{-\frac{3}{5}}$$

Typical  $C_n^2$  levels are between  $1 \times 10^{-16}$  (calm) to  $3 \times 10^{-13}$  (quite turbulent)

## The Coherent Doppler Lidar Equation, cont'd

The CNR equation can be written explicitly as

$$CNR(R) = \frac{\eta\beta T^2 c E_T}{h\nu B 2R^2} \frac{\pi D^2}{4} \left[ 1 + \left( \frac{\pi D^2}{4\lambda R} \right)^2 \left( 1 - \frac{R}{F} \right)^2 + \frac{D^2}{2\rho_o^2} \right]^{-1}$$

If the focus is at the range of interest, and if there is no turbulence, the CNR equation reduces to:

$$CNR(R) = \frac{\eta\beta T^2 c E_T}{h\nu B 2R^2} \frac{\pi D^2}{4}$$

# Next lecture...

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